

The Bose polaron- theory and experiments

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R. S. Christensen, J. Levinsen & GMB, PRL **115**, 160401 (2015)

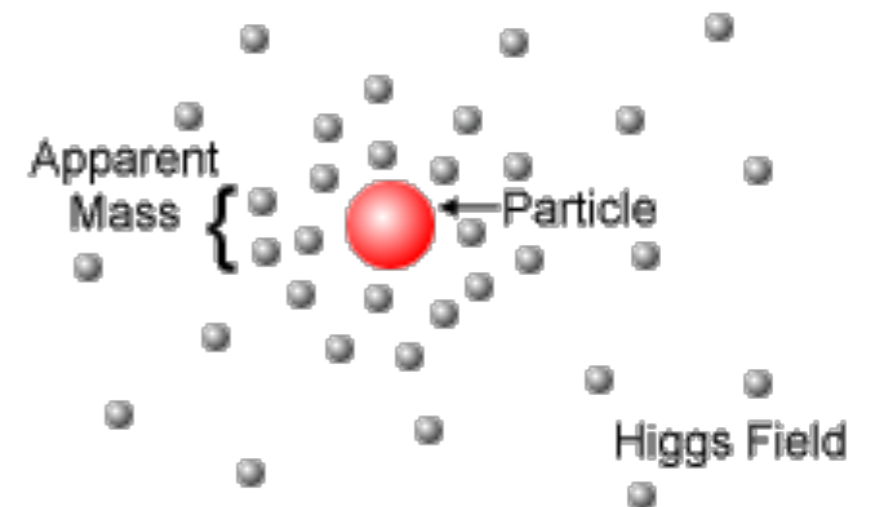
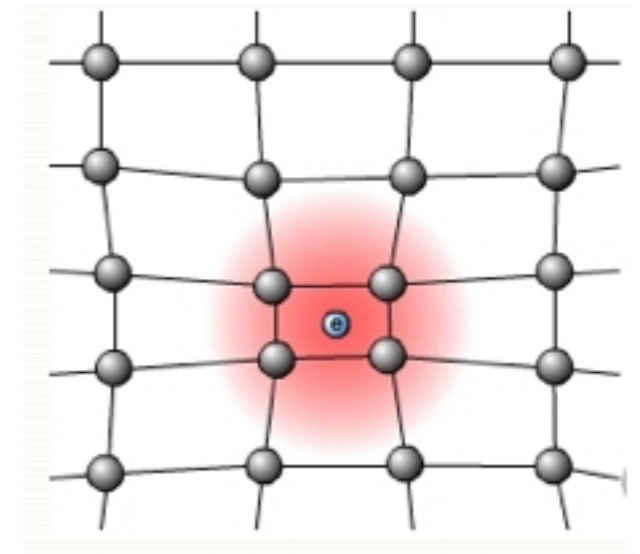
J. Levinsen, M. M. Parish & GMB, PRL **115**, 125302 (2015)

N. B. Jørgensen et al, arXiv:1604.07883

Bose Polaron

Mobile impurity interacting with *bosonic* reservoir

- Electrons coupled to phonons
- Helium mixtures
- High T_c superconductors
- Elementary particles coupled to the Higgs boson



Nice to have experimental realisation in cold atoms

Fermi polaron gave lots of new insights

1. Schirotzek *et al.*, Phys. Rev. Lett. **102**, 230402 (2009)
2. Kohstall *et al.*, Nature **485**, 615 (2012)
3. Koschourek *et al.*, Nature **485**, 619 (2012)



Very recently *two* independent experimental realisations of the Bose polaron:

- ① N. B. Jørgensen *et al.*, arXiv:1604.07883
- ② Ming-Guang Hu *et al.*, arXiv:1605.00729

This Talk

1. Theory

Good understanding, both at weak and strong coupling

2. Experiment

First observation of long lived Bose polaron using RF spectroscopy

People

Experiment



Jan Arlt



Nils
Jørgensen



Lars
Wacker



Kristoffer T.
Skalmstang



Aarhus University

Theory



Rasmus S.
Christensen



Jesper
Levinsen



Meera
Parish

Monash University



Theory

Mean-field: Astrakharchik & Pitaevskii, Phys. Rev. A **70**, 013608 (2004)
Cucchietti & Timmermans, Phys. Rev. Lett. **96**, 210401 (2006)
Kalas & Blume, Phys. Rev. A **73**, 043608 (2006)
Bruderer, Bao & Jaksch, Eu. Phys. Lett. **82**, 30004 (2008)

Fröhlich: Huang & Wan, Chin. Phys. Lett. **26**, 080302 (2009)
Tempere *et al.*, Phys. Rev. B **80**, 184504 (2009)
Castels & Wouters, Phys. Rev. A **90**, 043602 (2014)
Grust *et al.*, Sci. Rep. **5**, 12124 (2015)
Vlietinck *et al.*, New J. Phys. **17**, 033023 (2015)

Field theory: Rath & Schmidt, Phys. Rev. A **88**, 053632 (2013)

Variational: Li & Das Sarma, Phys. Rev. A **90**, 013618 (2014)
Schhaddilova, Schmidt, Grusdt & Demler, arXiv:1604.06469

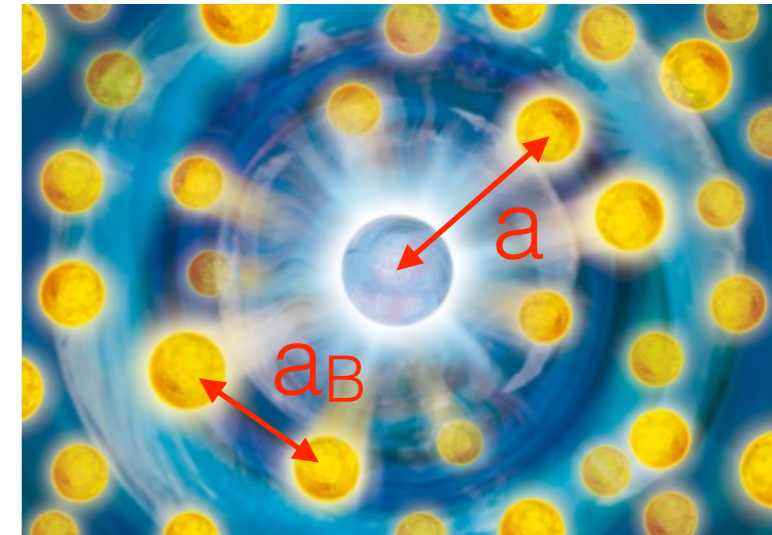
Perturbation theory

$$H = \sum_{\mathbf{k}} \epsilon_k^B a_{\mathbf{k}}^\dagger a_{\mathbf{k}} + \frac{1}{2\mathcal{V}} \sum_{\mathbf{k}, \mathbf{k}', \mathbf{q}} V_B(q) a_{\mathbf{k}+\mathbf{q}}^\dagger a_{\mathbf{k}'-\mathbf{q}}^\dagger a_{\mathbf{k}'} a_{\mathbf{k}} \quad \text{BEC}$$

$$+ \sum_{\mathbf{k}} \epsilon_k c_{\mathbf{k}}^\dagger c_{\mathbf{k}} + \frac{1}{\mathcal{V}} \sum_{\mathbf{k}, \mathbf{k}', \mathbf{q}} V(q) c_{\mathbf{k}+\mathbf{q}}^\dagger c_{\mathbf{k}'-\mathbf{q}}^\dagger a_{\mathbf{k}'} c_{\mathbf{k}}$$

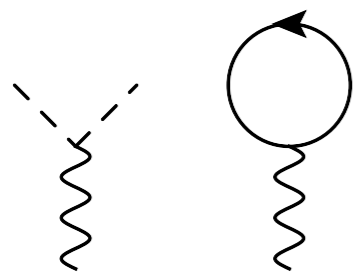
Impurity

Impurity-BEC interaction

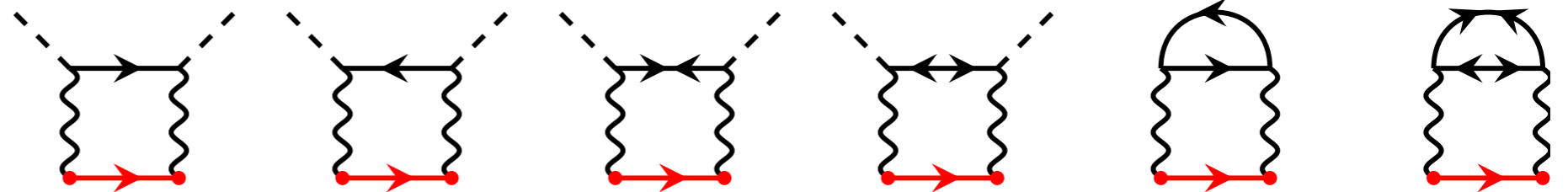


BEC weakly interacting $na_B^3 \ll 1 \Rightarrow$ Bogoliubov theory

Perturbation theory $V(q)$:



1. order



2. order

Replace $V(q) \rightarrow \mathcal{T}_v = 2\pi a/m_r$ in a consistent way

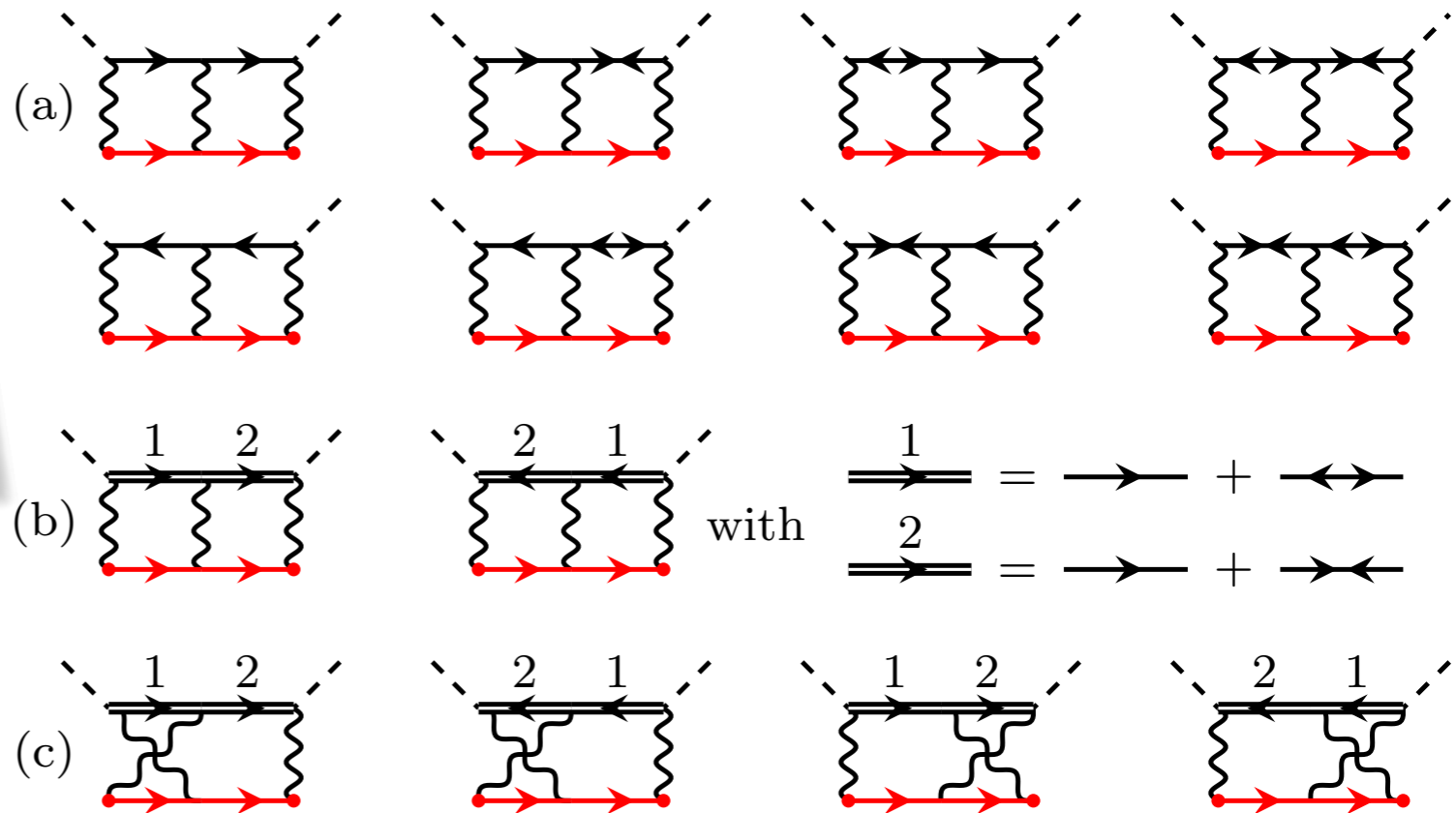
Diagrams like  comes from expanding

$$\mathcal{T}(p) = \frac{\mathcal{T}_v}{1 - \mathcal{T}_v \Pi_{11}(p)} = \mathcal{T}_v + \mathcal{T}_v^2 \Pi_{11}(p) + \dots$$

Self-energy in powers of a :

$$\Sigma(p, \omega) = \Sigma_1(p, \omega) + \Sigma_2(p, \omega) + \Sigma_3(p, \omega) + \dots$$

28 third order diagrams



Energy

$$E(0)$$

Same structure as
Lee-Huang-Yang +
Wu-Hugenholtz-Pines-Sawada

$$A(1) :$$

$$\alpha = n$$

$$= \max(a, a_B)$$

$$= 2\pi n \xi / m_r$$

a=a_B:

$$\frac{E}{N} = \frac{4\pi n a}{m} \left[1 + \frac{32}{3\sqrt{\pi}} (n a^3)^{1/2} + 4 \left(\frac{2}{3} \pi - \sqrt{3} \right) n a^3 \ln(n a^3) \right]$$

Weakly interacting BEC

$$\frac{E}{N} = \frac{2\pi n a}{m} \left[1 + \frac{128}{15\sqrt{\pi}} (n a^3)^{1/2} + 8 \left(\frac{4}{3} \pi - \sqrt{3} \right) n a^3 \ln(n a^3) \right]$$

Residue & Effective Mass

$$Z^{-1} = 1 + C(\alpha) \frac{a^2}{a_B \xi} + D(\alpha) \frac{a^3}{a_B \xi^2}$$

$$\frac{m^*}{m} = 1 + F(\alpha) \frac{a^2}{a_B \xi} + G(\alpha) \frac{a^3}{a_B \xi^2}$$

$$C(1) = 2\sqrt{2}/3\pi \quad D(1) \approx 0.64 \quad F(1) = 16\sqrt{2}/45\pi \quad G(1) \simeq 0.37$$

Condition for $Z \approx 1$: $\frac{a^2}{a_B \xi} \ll 1$

Breaks down for ideal BEC

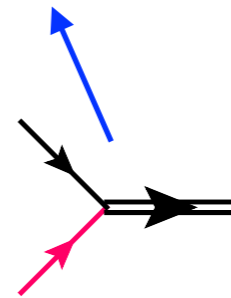
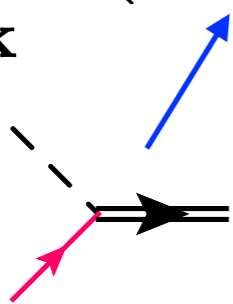
Variational Theory

Multichannel model

$$\hat{H} = \sum_{\mathbf{k}} \left[E_{\mathbf{k}} \beta_{\mathbf{k}}^{\dagger} \beta_{\mathbf{k}} + \epsilon_{\mathbf{k}} c_{\mathbf{k}}^{\dagger} c_{\mathbf{k}} + (\epsilon_{\mathbf{k}}^{\text{d}} + \nu_0) d_{\mathbf{k}}^{\dagger} d_{\mathbf{k}} \right]$$

Bog. modes Impurity Molecule

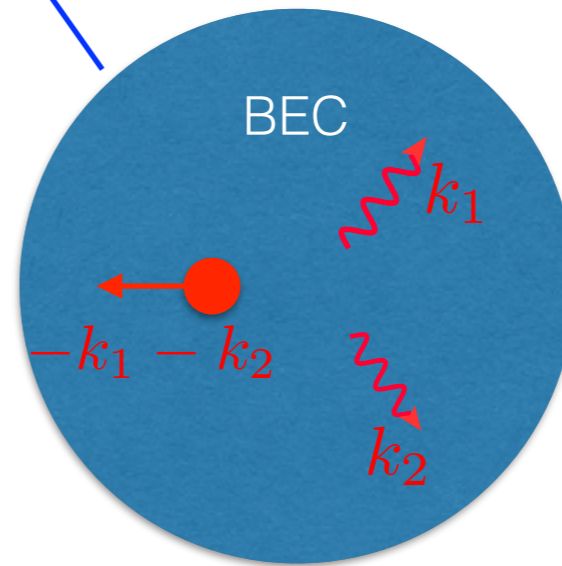
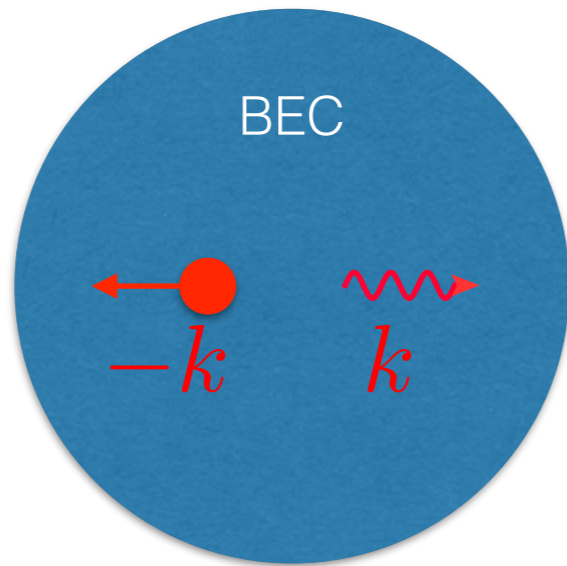
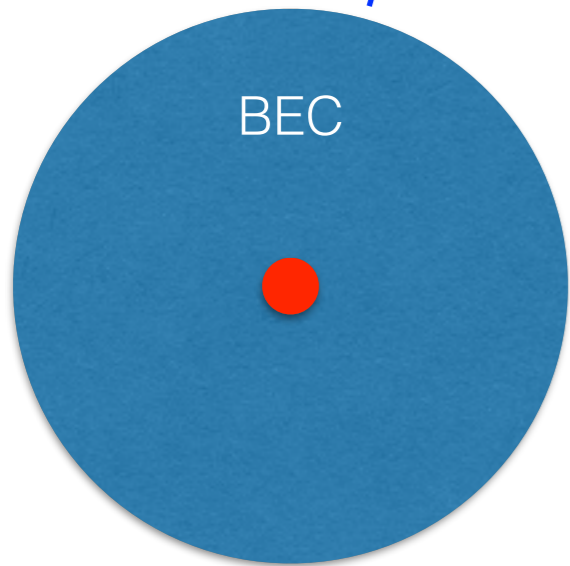
$$+ g \sqrt{n_0} \sum_{\mathbf{k}} \left(d_{\mathbf{k}}^{\dagger} c_{\mathbf{k}} + h.c. \right) + g \sum_{\mathbf{k}, \mathbf{q}} \left(d_{\mathbf{q}}^{\dagger} c_{\mathbf{q}-\mathbf{k}} b_{\mathbf{k}} + h.c. \right)$$



Introduces effective range r_0
Regularises 3-body problem

Variational wave function

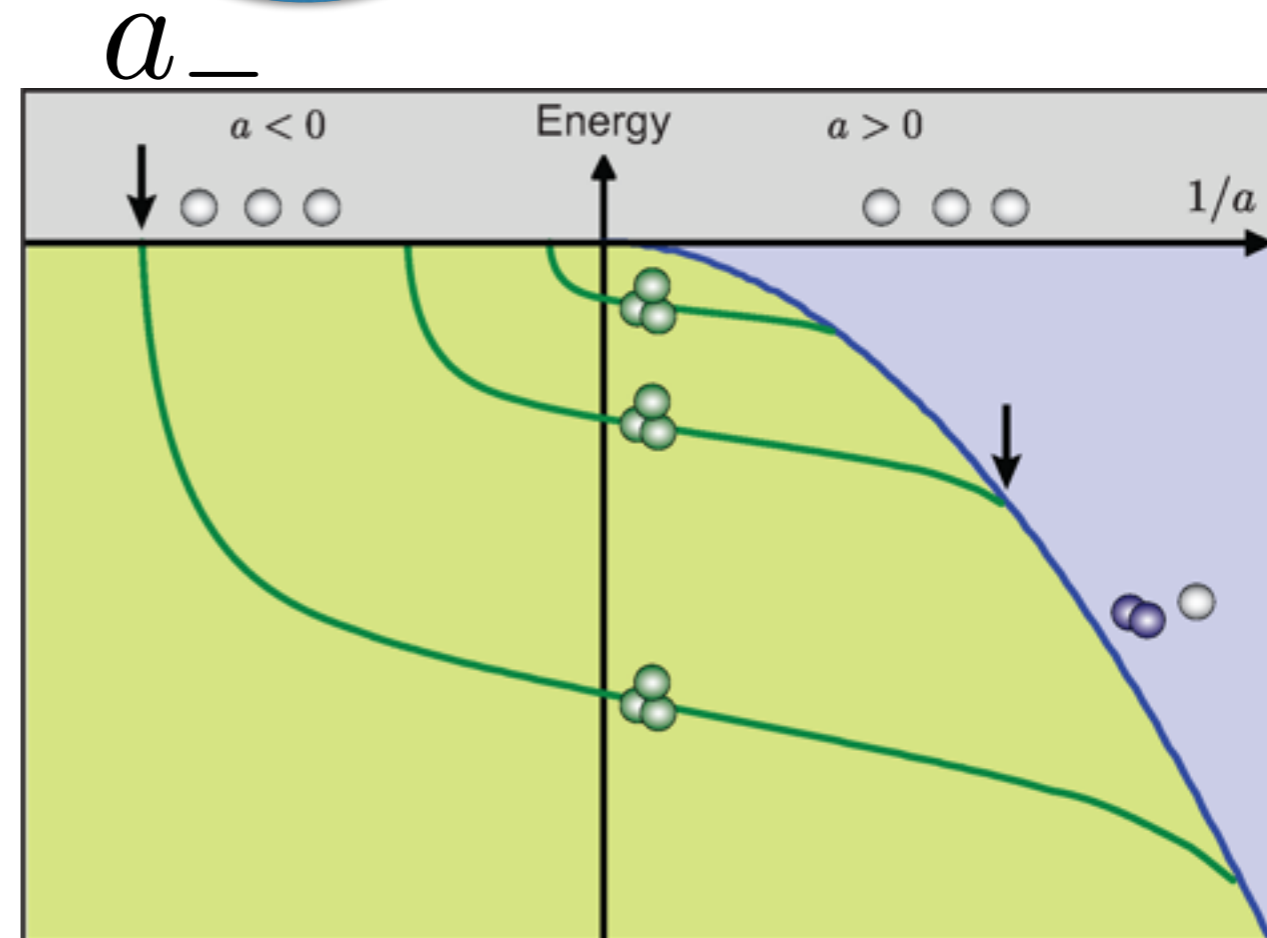
$$|\psi\rangle = \left(\alpha_0 c_0^\dagger + \sum_{\mathbf{k}} \alpha_{\mathbf{k}} c_{-\mathbf{k}}^\dagger \beta_{\mathbf{k}}^\dagger + \frac{1}{2} \sum_{\mathbf{k}_1 \mathbf{k}_2} \alpha_{\mathbf{k}_1 \mathbf{k}_2} c_{-\mathbf{k}_1 - \mathbf{k}_2}^\dagger \beta_{\mathbf{k}_1}^\dagger \beta_{\mathbf{k}_2}^\dagger + \gamma_0 d_0^\dagger + \sum_{\mathbf{k}} \gamma_{\mathbf{k}} d_{-\mathbf{k}}^\dagger \beta_{\mathbf{k}}^\dagger \right) |\text{BEC}\rangle$$

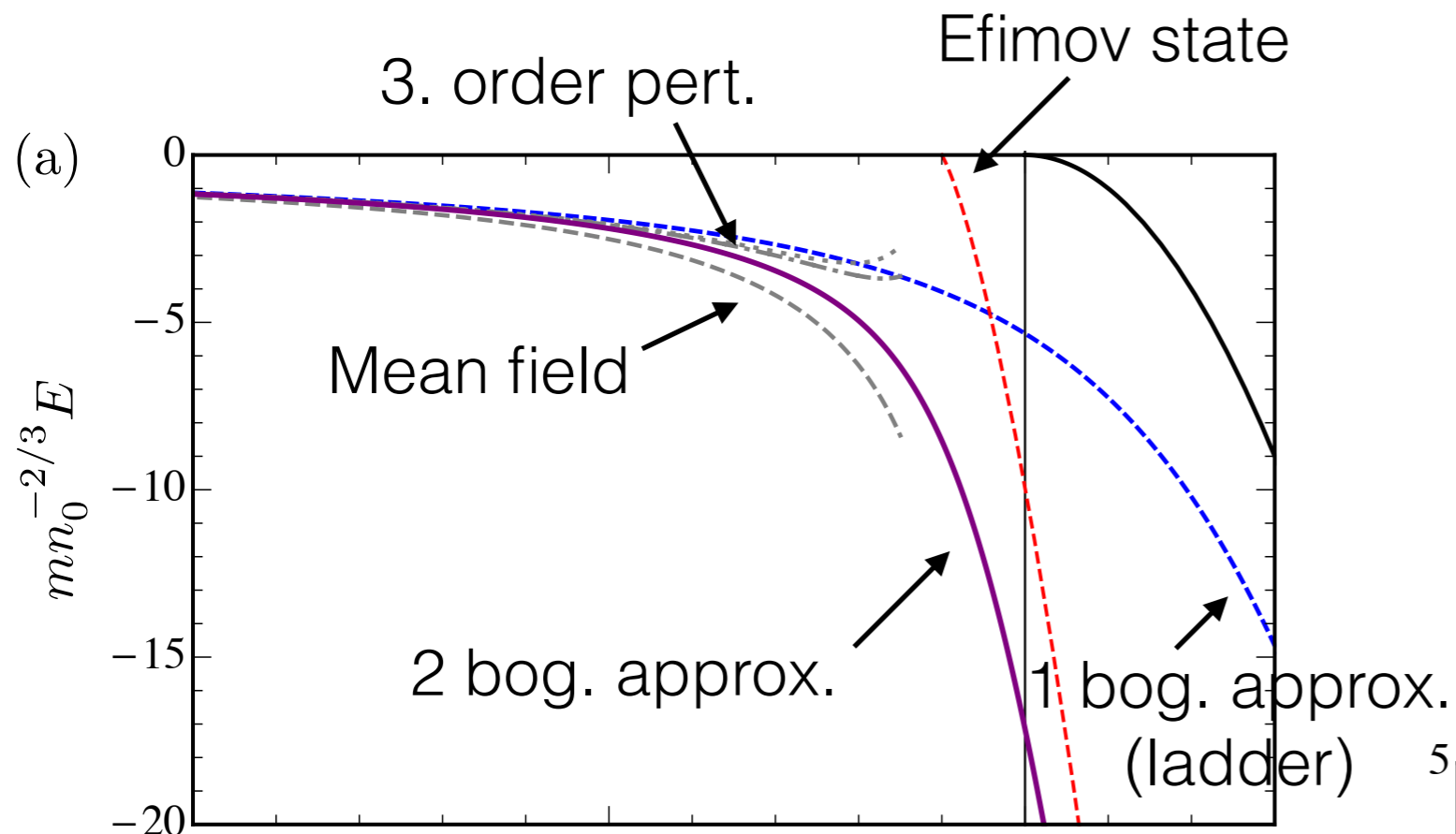


Recovers Efimov spectrum
for 1+2 bosons for $n_0 \rightarrow 0$

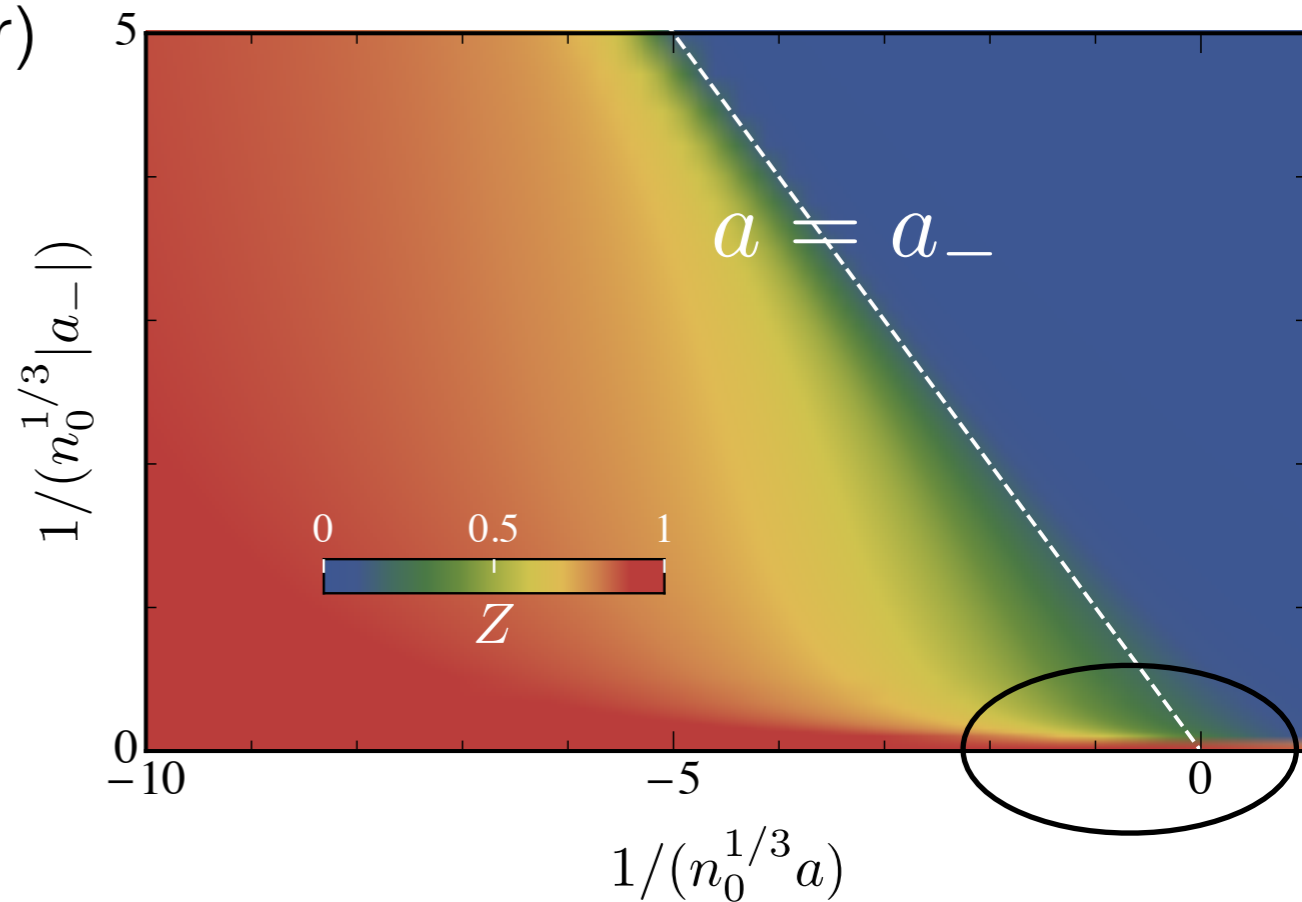
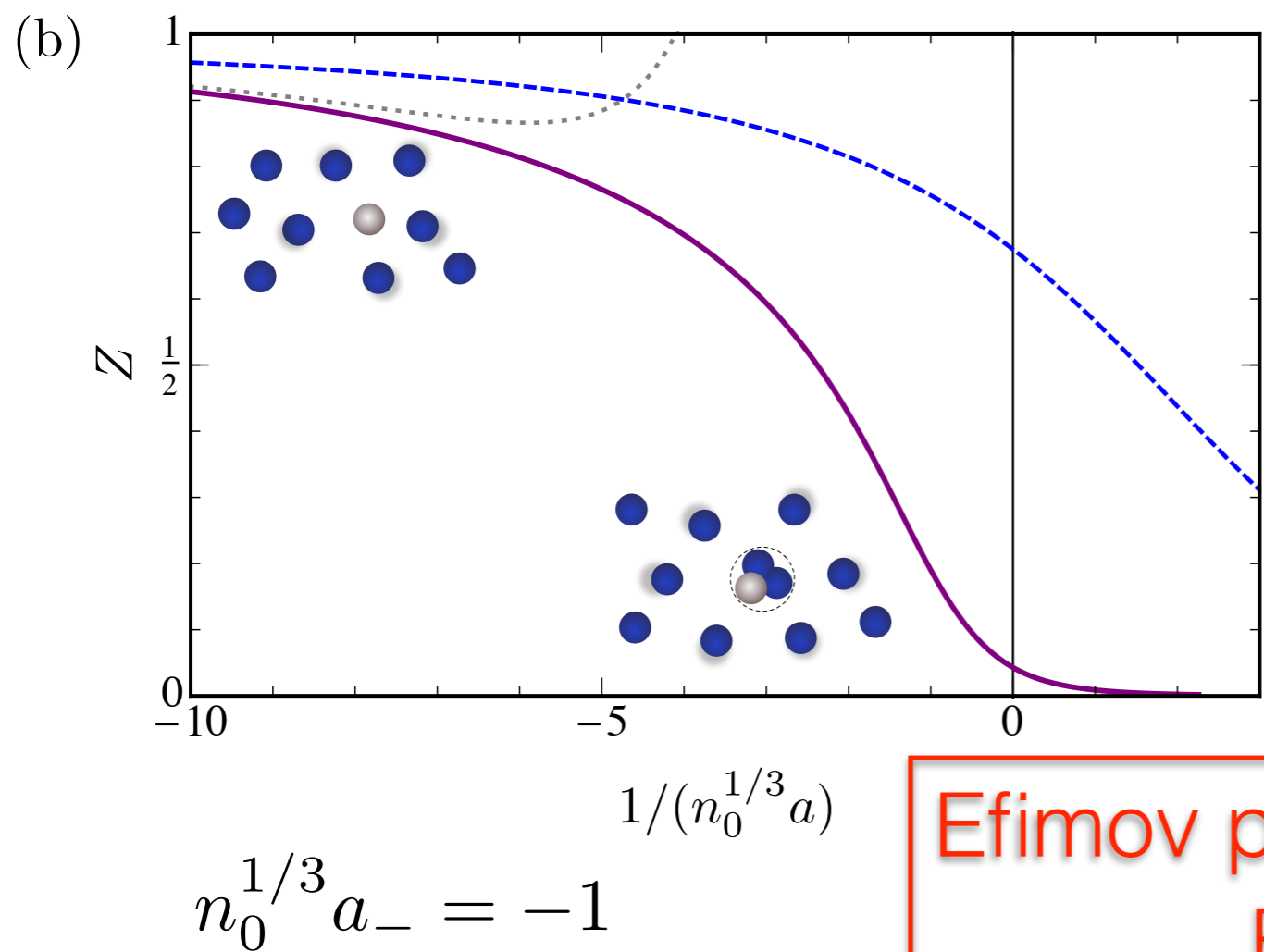
$$a_- \simeq 9000 r_0$$

Can always keep $|a/r_0| \gg 1$
Even when a_- large





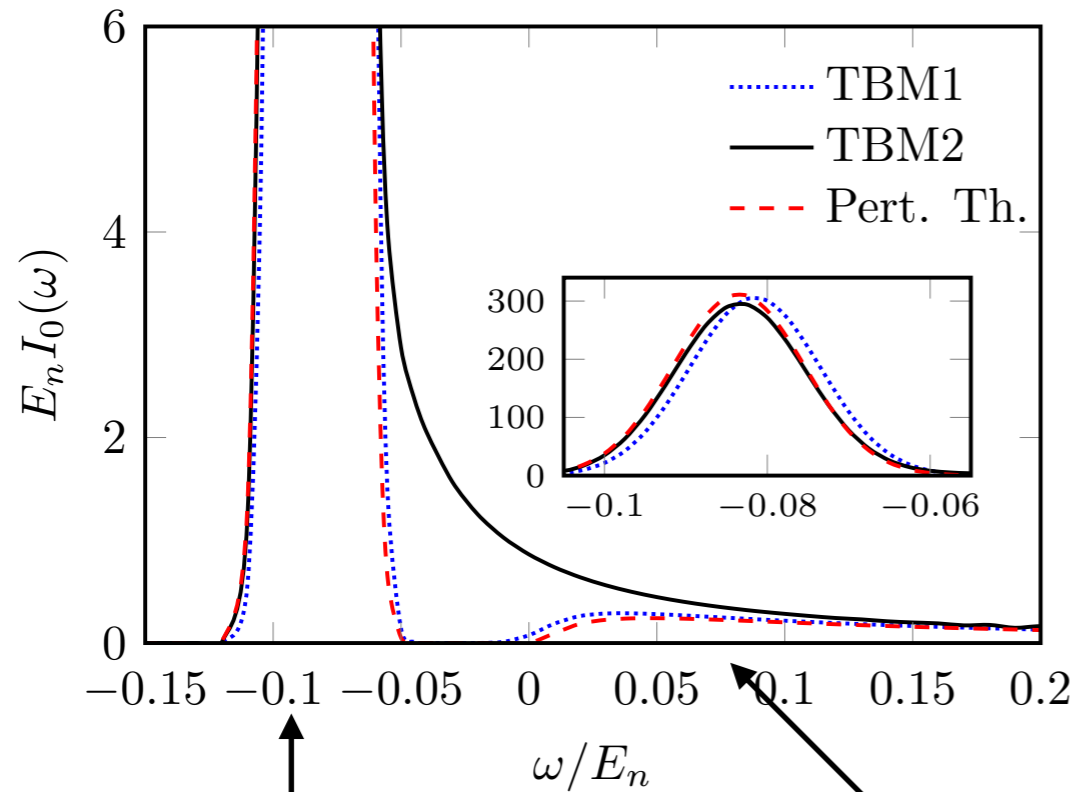
Avoided crossing
with Efimov state.
Residue small close
to unitarity



Efimov physics suppressed for $a \gg n_0^{-1/3}$
Efimov state "too large"

Spectral functions

$$1/k_n a = -5$$



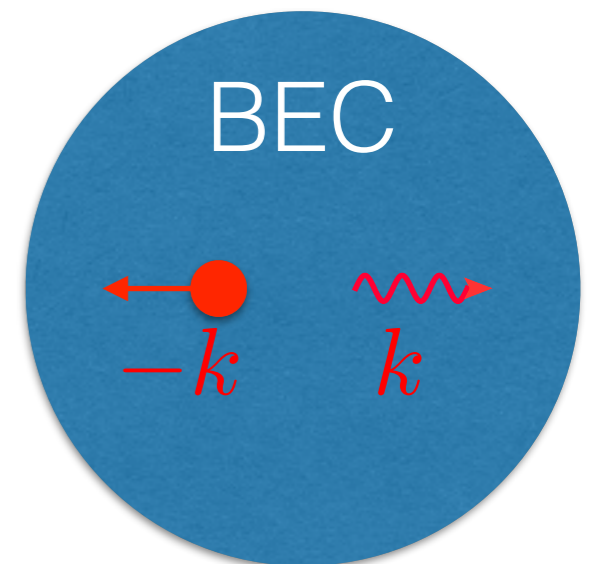
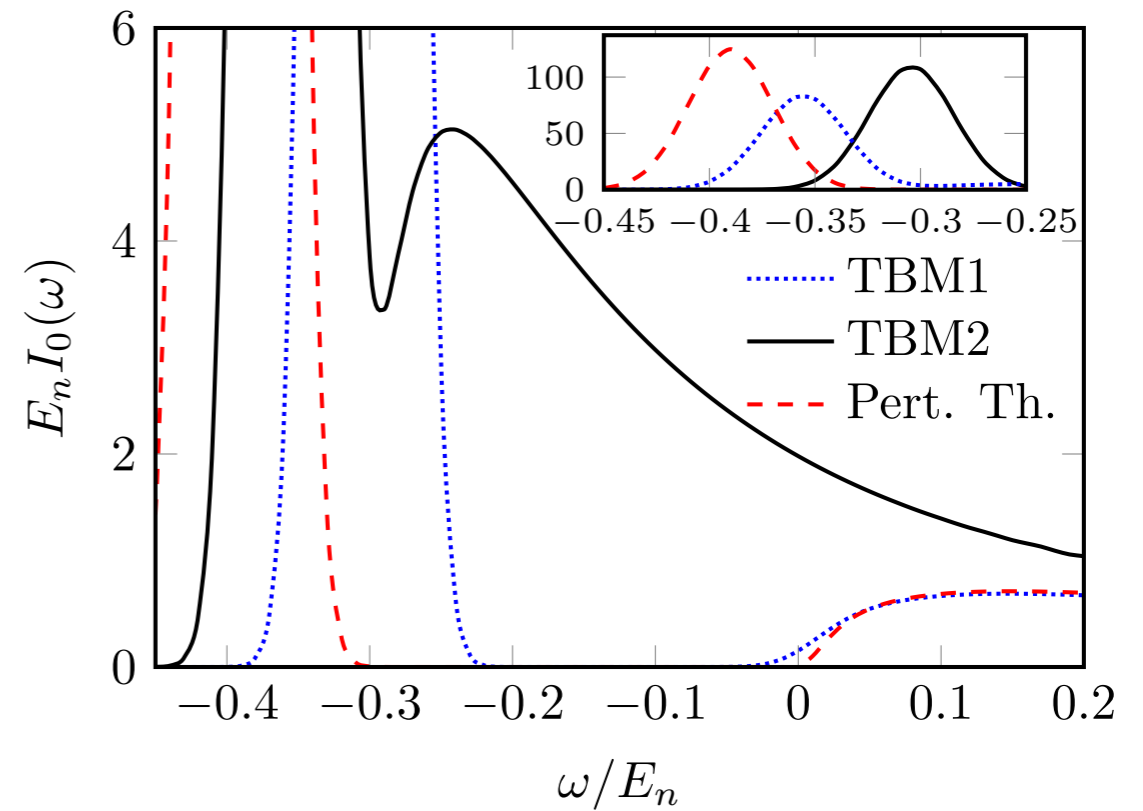
Polaron peak

Many-body continuum

Weak coupling: Variational theory agrees with pert. theory

Strong coupling: Pert. theory breaks down. Many-body continuum significant

$$1/k_n a = -1$$



Theory - bottom lines

- Analytical perturbation theory to 3.order in a/ξ
- Polaron well-defined for weak coupling
- Strong coupling: Variational ansatz including 3-body Efimov correlations
- Significant many-body continuum for strong coupling
- Impurity atoms in BEC not the Fröhlich model

Experiment bottom lines



Jan Arlt

- First realisation of the Bose polaron (See also JILA group)
- Well-defined polaron both for repulsive and attractive interaction
- Many-body continuum dominates at strong coupling
- Excellent agreement with theory
- 3-body decay has no significant effects

Earlier experiments

Impurity in thermal
bose gas:

Spethman *et al.* Phys. Rev. Lett. **109**, 235301 (2012)

Charged or fixed
impurities in BEC:

Zipkes *et al.*, Nature **464**, 388 (2010)

Schmid *et al.*, Phys. Rev. Lett. **105**, 133202 (2010)

Balewski *et al.*, Nature **502**, 664 (2013)

Scelle *et al.*, Phys. Rev. Lett. **111**, 070401 (2013)

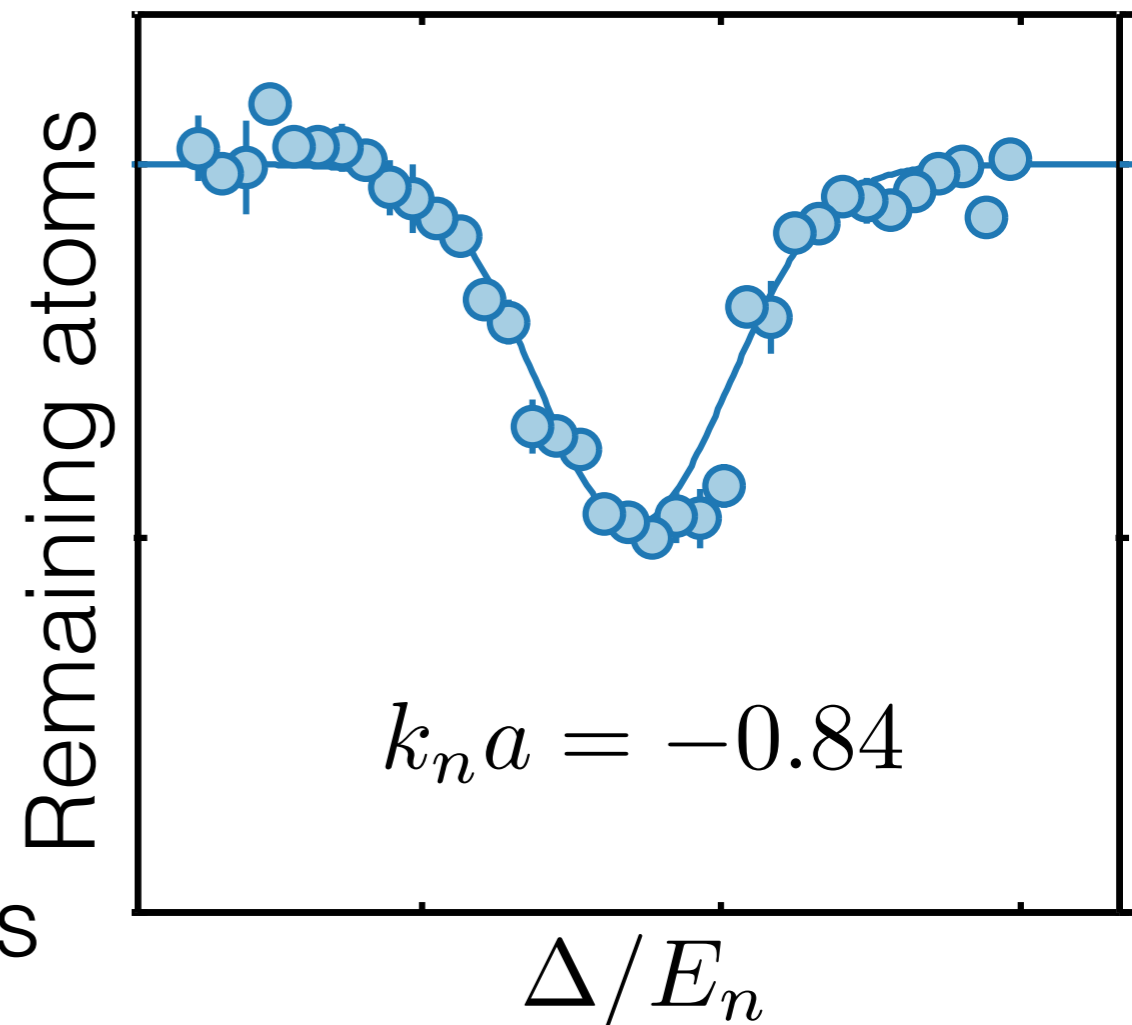
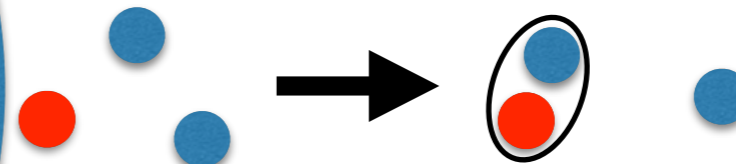
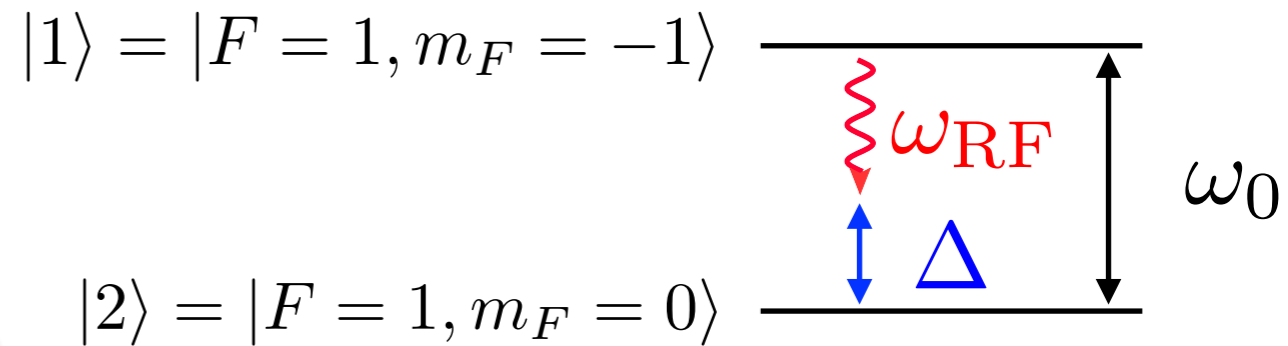
Impurities in lattice: Ospelkaus *et al.*, Phys. Rev. Lett. **96**, 180403 (2006)

Magnons:

Marti *et al.*, Phys. Rev. Lett. **113**, 155302 (2014)

Experimental procedure

BEC of ^{39}K in $|1\rangle$
 RF flip $\leq 10\%$ to $|2\rangle$
 Wait for a while
 TOF
 } 3-body loss
 Count # $|1\rangle$ remaining as f_n
 of detuning $\Delta = \omega_0 - \omega_{\text{RF}}$

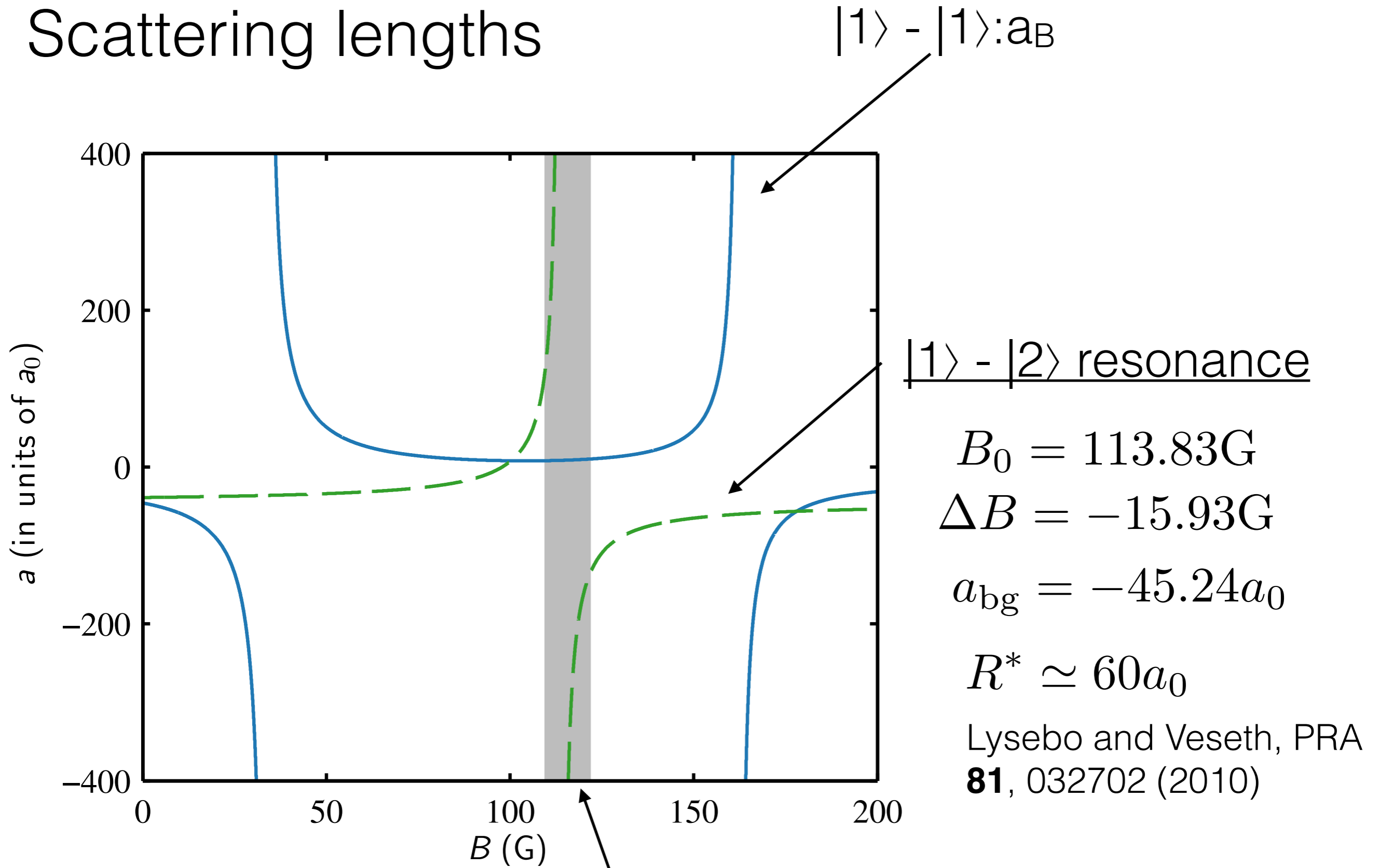


Independent of wait time \Rightarrow

lose 100% of $|2\rangle$ atoms \Rightarrow

lost $|1\rangle$ atoms = $3 \times$ created $|2\rangle$ atoms

Scattering lengths



Experiment takes place here. $a_B \approx 9a_0$

Advantages of RF flipping out of BEC

- ❶ Perfect spatial overlap between impurities and BEC
- ❷ Selectively probe only $k=0$ polarons
- ❸ Simple theoretical interpretation

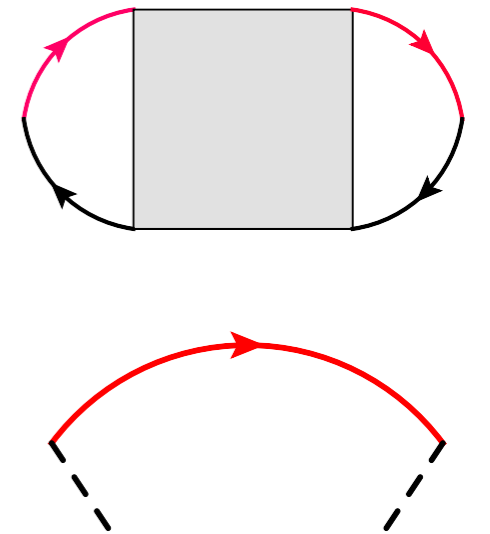
$$\dot{N}_2 = -2\Omega^2 \text{Im}D(\omega)$$

$$D(t-t') = -i\theta(t-t') \langle [\sum_{\mathbf{k}} a_{\mathbf{k}1}^\dagger(t) a_{\mathbf{k}2}(t), \sum_{\mathbf{k}'} a_{\mathbf{k}'2}^\dagger(t') a_{\mathbf{k}'1}(t')] \rangle$$

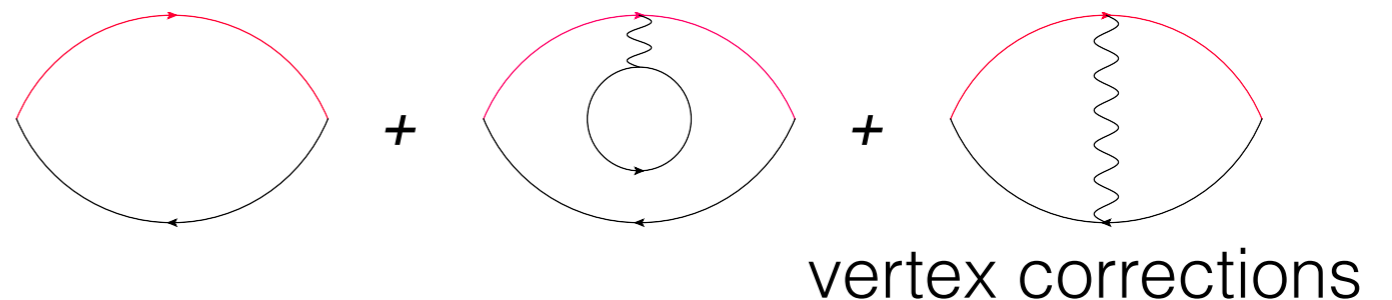
Bogoliubov theory: $D(\omega) = n_0 G_2(\mathbf{k} = \mathbf{0}, \omega)$

RF probes $k=0$ impurity spectral function:

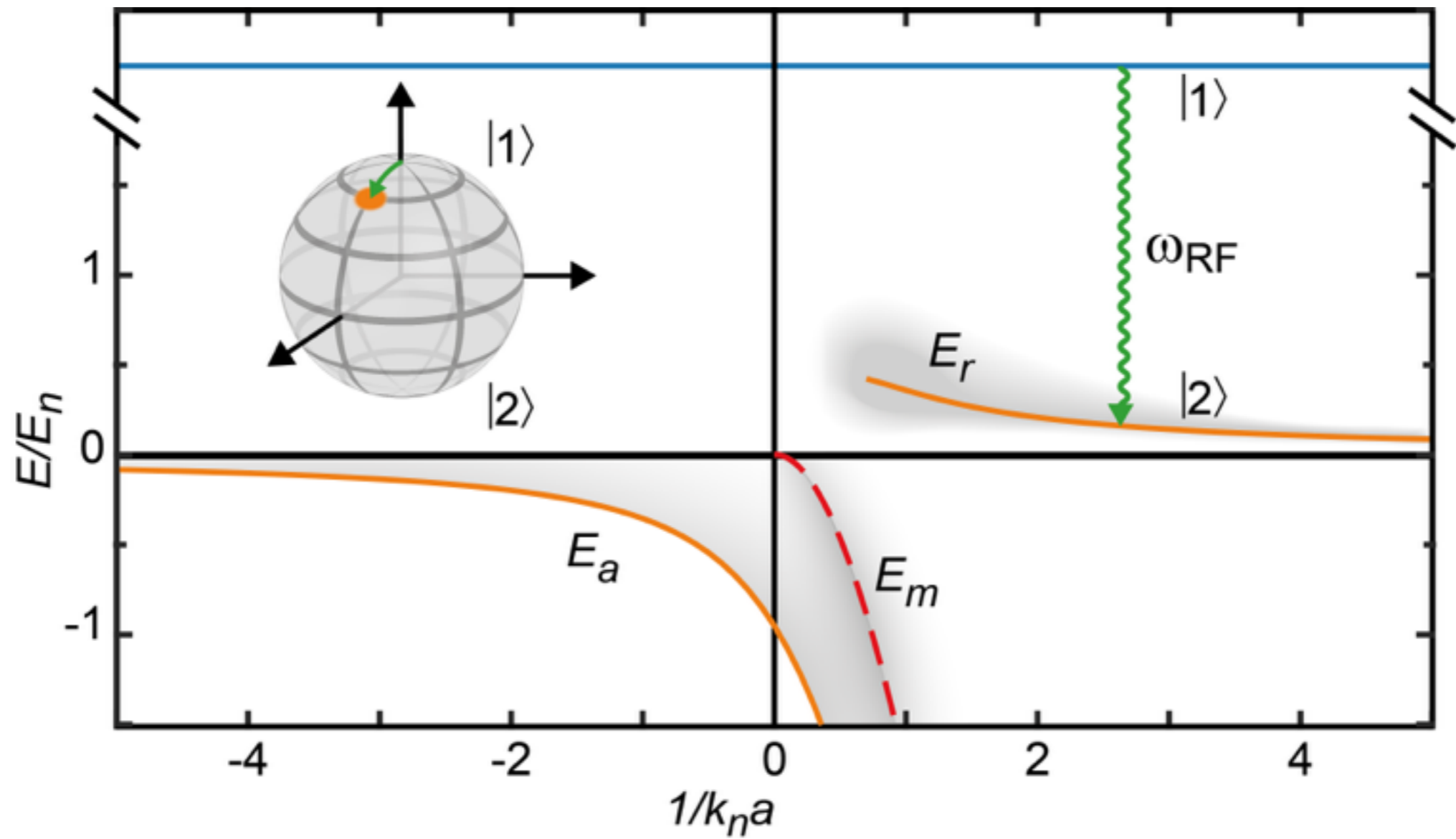
$$\dot{N}_2 \propto A(\mathbf{k} = \mathbf{0}, \omega) = -2\text{Im}G_2(\mathbf{k} = \mathbf{0}, \omega)$$



Contrast with Fermi gas or thermal Bose gas



Generic Physics

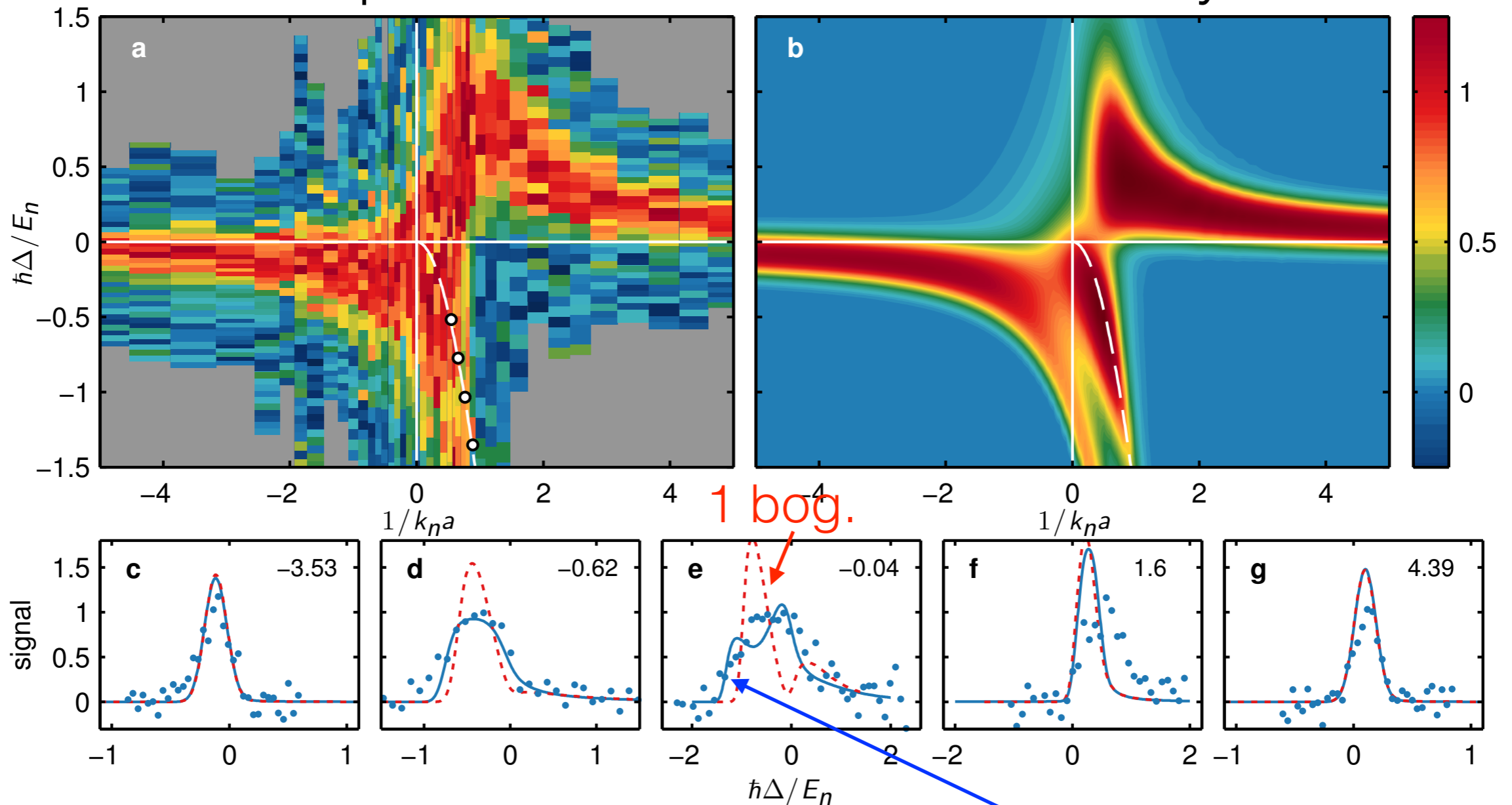


$$k_n = (6\pi^2 n)^{1/3}$$

$$E_n = \frac{k_n^2}{2m}$$

Experiment

Theory



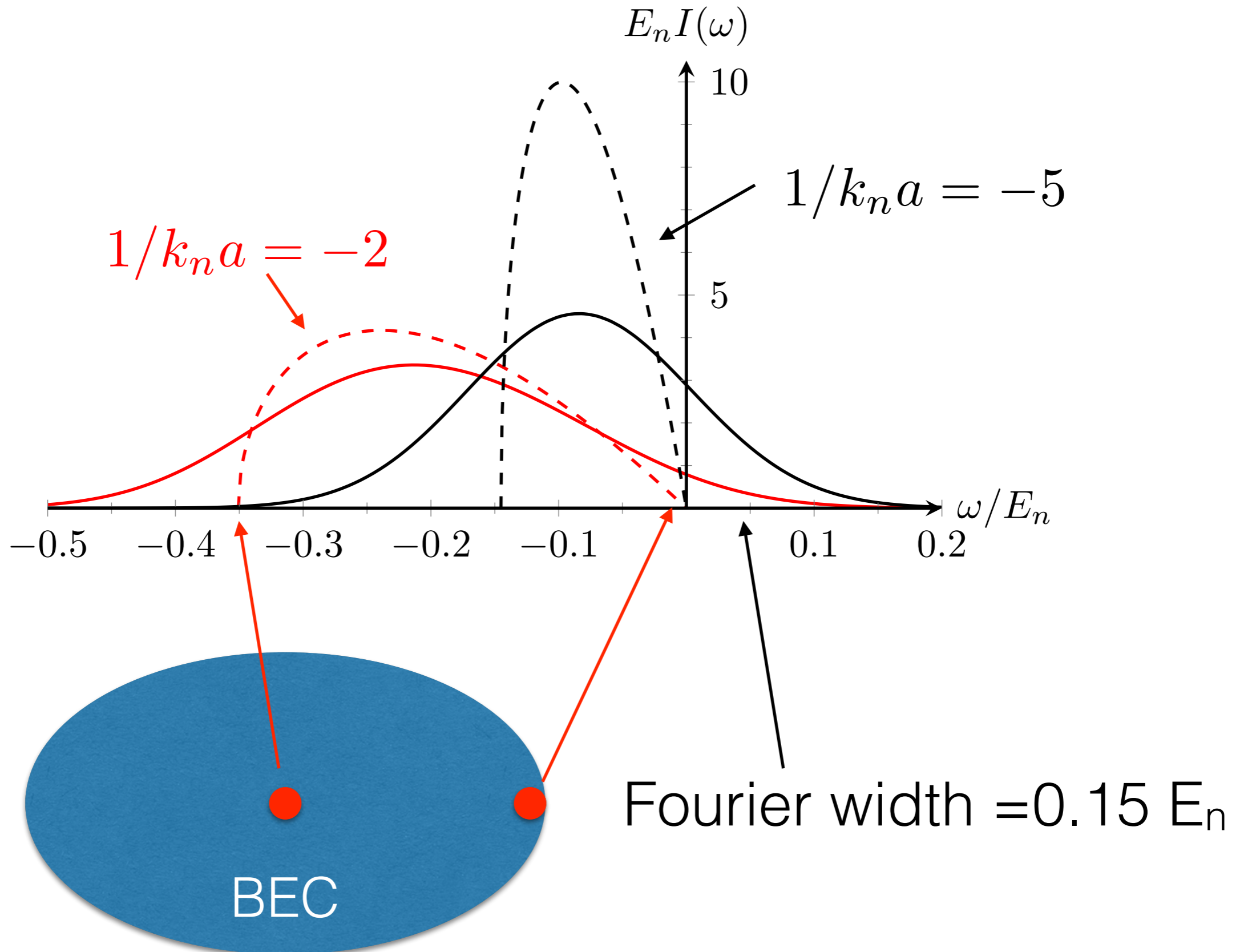
★ Clear shift away from ω_0

★ Excellent agreement between experiment and 2 bog. theory (trap averaging important!)

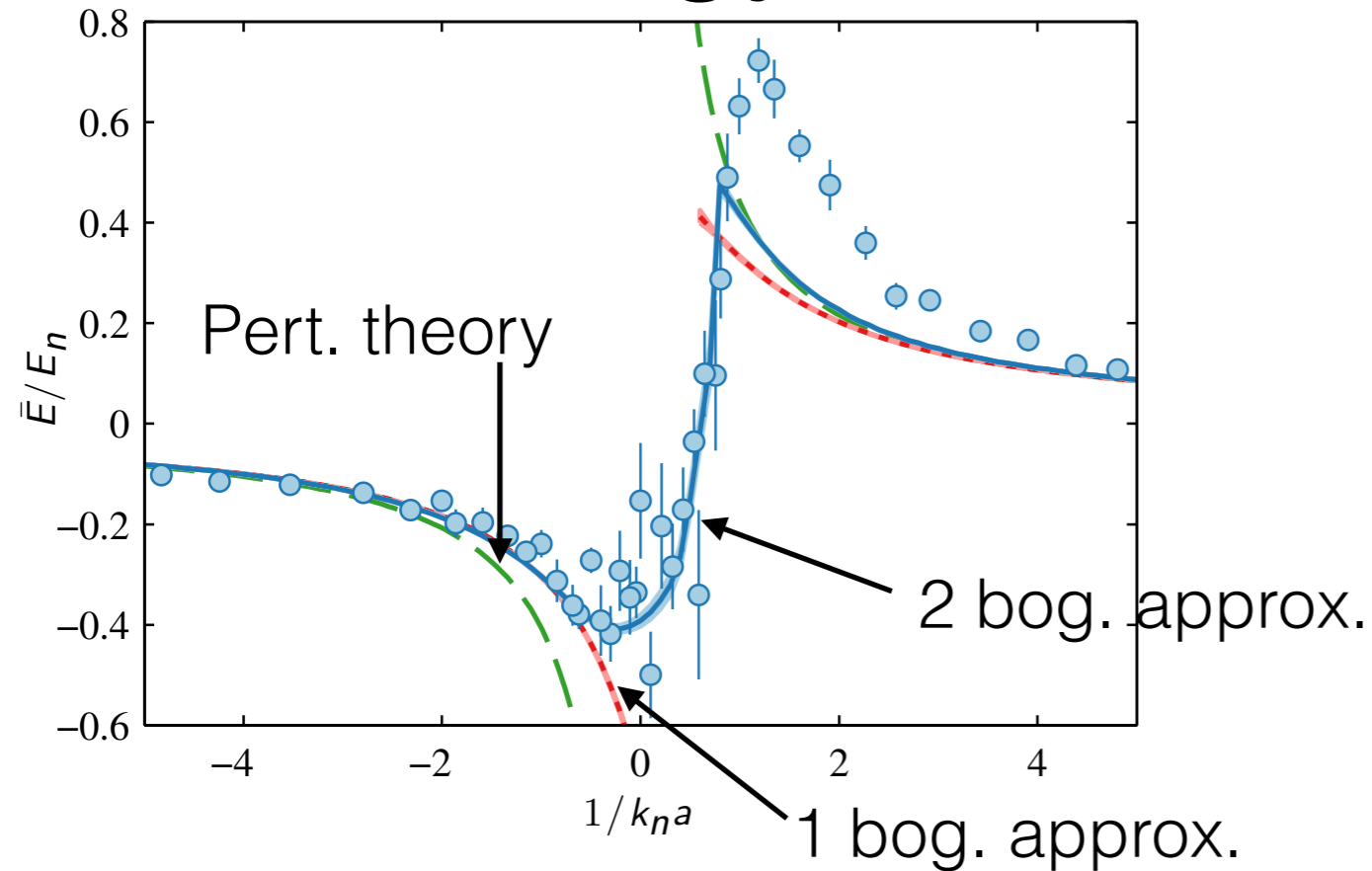
★ Well-defined polaron for weak coupling

★ Many-body continuum dominates for strong coupling

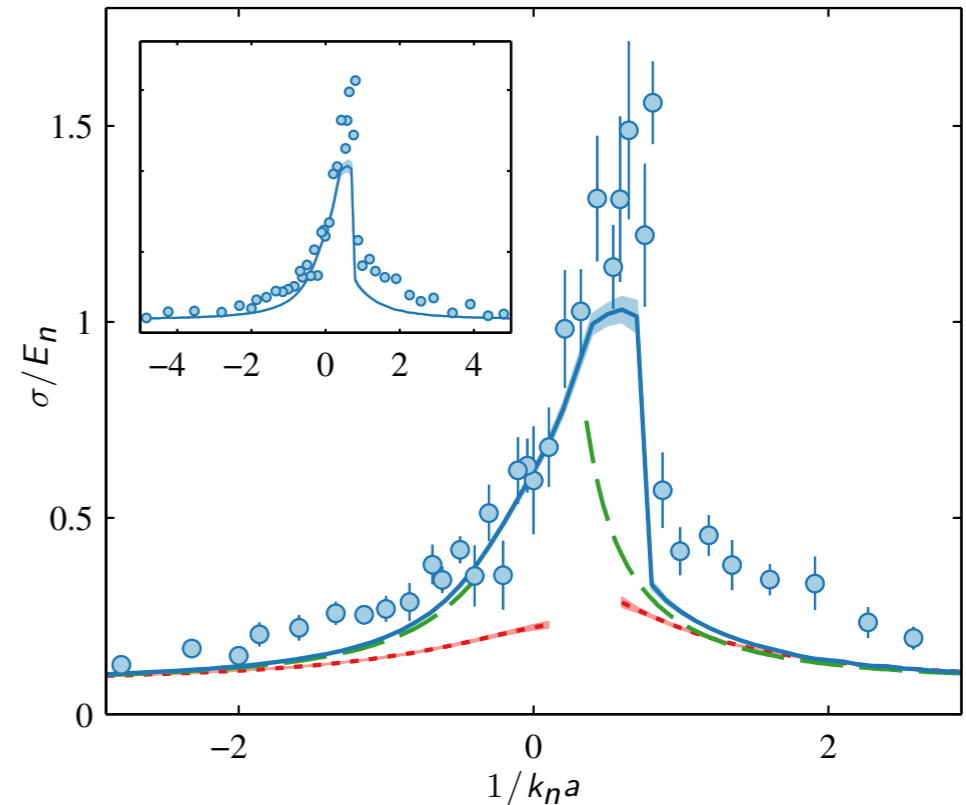
Trap averaging & Fourier broadening



Energy



Width



★ Remarkable agreement between experiment and theory
(some problems at strong repulsion)

★ Pert. theory explains data for weak coupling \Rightarrow
well defined polaron

★ 3-body decay not needed
to explain width

$$\Gamma \propto n_0^2 a^4 \text{ weak coupling}$$

$$\Gamma \propto E_n \text{ unitarity}$$

Makotyn *et al.*, Nat. Phys. **10**, 116 (2014)

Conclusions

- ➊ Good theoretical understanding of Bose polaron both for weak and for strong coupling
- ➋ Experimental observation of Bose polaron for the first time

R. S. Christensen, J. Levinsen & GMB, PRL **115**, 160401 (2015)

J. Levinsen, M. M. Parish & GMB, PRL **115**, 125302 (2015)

N. B. Jørgensen et al, arXiv:1604.07883



EFB23

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DEPARTMENT OF MATHEMATICS, AARHUS UNIVERSITY, DENMARK

8TH - 12TH AUGUST 2016

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Chen Ji, ECT* Trento

Chris Greene, Purdue University

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Elzbieta Stephan, University of Silesia

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Francesca Sammarruca, University of Idaho

Frank Deuretzbacher, ITP University of Hannover

Laura Marcucci, Pisa University

Lucas Platter, University of Tennessee

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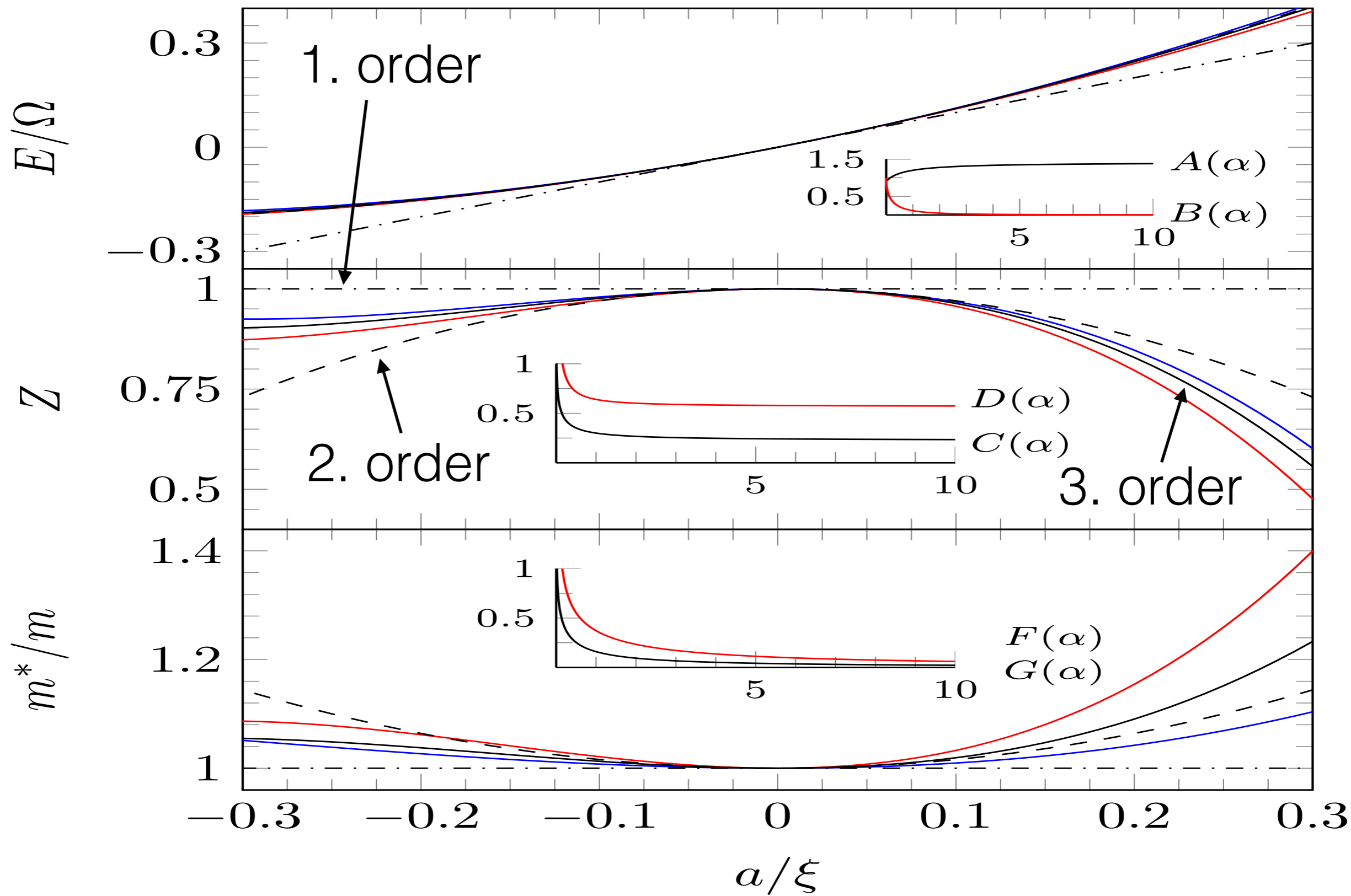
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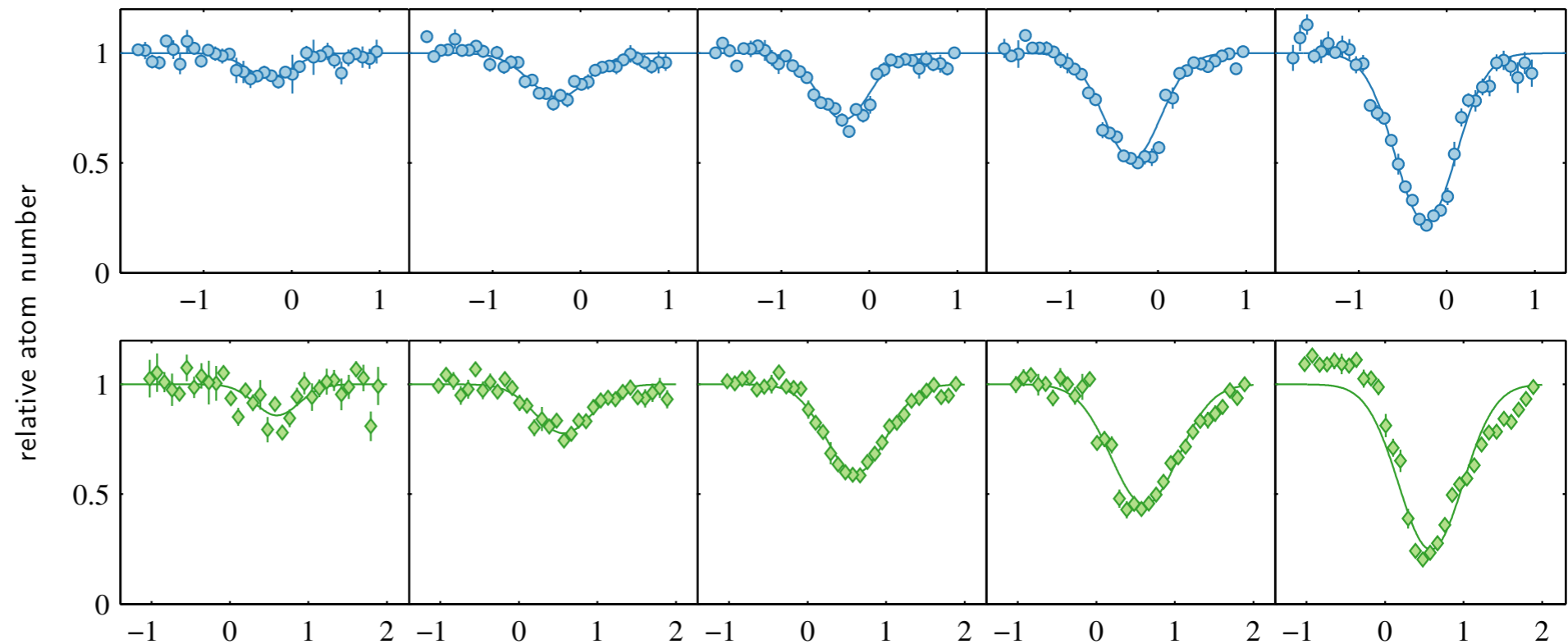
$$a_B/\xi = 0.1$$

Linear Response regime

Increasing RF power

$$1/k_n a = -0.84$$

$$1/k_n a = 1.6$$



Position and width stable

