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## Local magnetic fields in antiferromagnetic Bi<sub>2</sub>CuO<sub>4</sub>: as seen from <sup>63,65</sup>Cu and <sup>209</sup>Bi nuclear resonance

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## Abstract

A complex spin-echo spectrum of  ${}^{63,65}$ Cu and  ${}^{209}$ Bi has been observed in antiferromagnetic Bi<sub>2</sub>CuO<sub>4</sub> at 4.2 K. The spectrum consists of  ${}^{63,65}$ Cu AFNR and  ${}^{209}$ Bi NQR. The Cu AFNR is characterized by the internal field  $H_{int} = 99$  kOe at Cu nuclei and by the  ${}^{63}v_Q = 10$  MHz. It is argued that the low-limit of the internal field at the  ${}^{209}$ Bi nuclei is 4.3 kOe. The occurrence of  $H_{hf}$  at  ${}^{209}$ Bi nuclei shows that Bi $^{3+}$  anions take part in the superexchange interaction between Cu atoms along the Cu–O–Bi–O–Cu bonds. © 2000 Published by Elsevier Science B.V. All rights reserved.

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The structure of  $Bi_2CuO_4$  does not allow an ordinary Cu–O–Cu superexchange interaction between Cu<sup>2+</sup> cations. The AF order is maintained by strongly anisotropic inter- and intra-sublattice interactions involving at least several anions by two possible ways: Cu–O–Bi–O–Cu [1] or Cu–O–O–Cu [2]. There is still no understanding of the role of Bi ions in these interactions. Here we present spin-echo studies in pure AF samples of  $Bi_2CuO_4$ .

Cu NQR spectra were obtained using a point-by-point technique at 4.2 K. In order to separate the  ${}^{63,65}$ Cu part of the observed complex spectrum from  ${}^{209}$ Bi we measured the zero-field spectrum of Bi<sub>2</sub>CuO<sub>4</sub> enriched by the  ${}^{65}$ Cu isotope (Figs. 1–3). This allows to conclude, that all the observed resonance lines below 80 MHz belong to  ${}^{209}$ Bi nuclei.

We have assigned the upper frequency triplet with the intense central line at 112.5 MHz in Fig. 3 to a quadrupole splitted spectrum of <sup>63</sup>Cu isotope. Using the experimental frequencies of the left- and right-side peaks

(107.0(1) and 117.0(1) MHz) of the triplet, the secondorder perturbation treatment for  $\eta = 0$  and  $\theta = 90^{\circ}$ yields  ${}^{63}v_{\rm L} = 112.0(1)$  MHz and  ${}^{63}v_{Q} = 10.0(1)$  MHz. The calculated frequency for the central line, 112.2 MHz, lies within the limits of experimental error of peak position.

The obtained Larmor frequency  $^{63}\nu_{\rm L} = 112.0(1)$  MHz corresponds to a hyperfine field of 99.0(1) kOe. Using the hyperfine coupling constant for Cu<sup>2+</sup> equal to 120 kOe/ $\mu_{\rm B}$  and ignoring the dipolar contribution, we obtain a magnetic moment of  $0.82\mu_{\rm B}$  which is close to the value of  $0.85\mu_{\rm B}$  deduced from neutron diffractometry [3].

From the principal axes directions of the electric field gradient tensor it is evident that in the elementary cell of  $Bi_2CuO_4$  there are four pairs of crystallographic sites of Bi atoms which become inequivalent when a magnetic field appears at Bi sites. Such a magnetic inequivalency might be the reason for the abundance of resonance lines observed in  $Bi_2CuO_4$  (Figs. 1–3), which considerably hampers the interpretation of the spectrum.

In the presence of a high internal magnetic field  $H_{\text{int}}$  and the asymmetry parameter  $\eta$  close to unity, the Bi nuclei could exhibit Zeeman transitions ( $-m \Leftrightarrow m$ ) with

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Fig. 1. Low-frequency part of the zero-field spectra of  $Bi_2CuO_4$ . Ceramic sample enriched by <sup>65</sup>Cu. Inset: the low-frequency line of this spectrum in the same frequency scaling.



Fig. 2. Middle-frequency part of the zero-field spectra of  $Bi_2CuO_4$ . Ceramic sample enriched by <sup>65</sup>Cu.

the highest probability for  $(-\frac{1}{2} \Leftrightarrow \frac{1}{2})$  [4]. Since the low-frequency spectrum cannot fit Bi NQR, we assume that the resonance line at 26.35 MHz corresponds to the  $(-\frac{1}{2} \Leftrightarrow \frac{1}{2})$  transition. We estimate the minimum value of  $H_{\text{int}}$  necessary to create this transition to be  $H_{\text{int}}(\text{min}) = 4.3 \text{ kOe}.$ 



Fig. 3. High-frequency part of the zero-field spectra of different  $Bi_2CuO_4$  samples: (a,b) different single crystals and (c) ceramic sample enriched by  $^{65}Cu$ .

This value is almost one order of magnitude larger than  $H_{dip}$ , which supports our assumption about a comparatively large  $H_{hf}$  contribution to  $H_{int}$  at Bi atoms.

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