

Laser-induced Femtosecond Spin Dynamics in Metallic Multilayers

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13.08.2009



1 Introduction

- Spin-polarization
- Spin-transfer

2 Experimental Realisation

- Optical setup
- First tests

3 Experimental Results

- Measured data

4 Outlook

- Optimization of spin-emission
- Investigation of hot electron transport
- Spin-transfer by hot carriers

5 Summary

- Summary
- Thank You

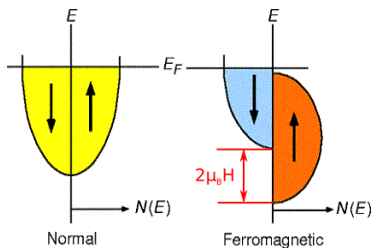


Spin-polarization in Fe

- energy spin splitting of valence band (different energies for *spin up* and *spin down*)
- therefore spin polarization ($N^{up} \neq N^{down}$)
- macroscopic magnetization due to spin polarization:

$$M = \mu_B \cdot (N^{up} - N^{down})$$

where $N^{up} - N^{down} \propto 2\mu_B H$

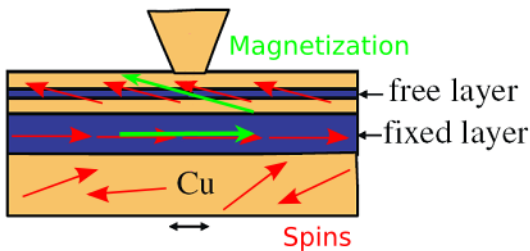


G.A. Prinz, *Science* 282 (1998) 1160



Spin-transfer by current

(Ralph, Stiles, JMMM 320 1190 (2008))

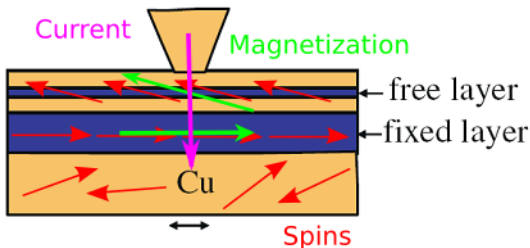


1. magnetization \vec{M} in free and fixed layer is non-parallel



Spin-transfer by current

(Ralph, Stiles, JMMM 320 1190 (2008))

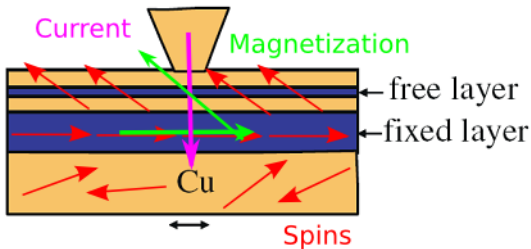


2. apply an external bias to induce a current from the bottom Cu-layer; the current is spin-polarized by the fixed layer



Spin-transfer by current

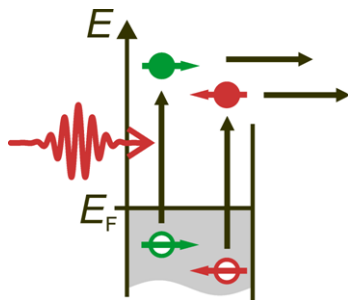
(Ralph, Stiles, JMMM 320 1190 (2008))



3. the spin-polarized current changes \vec{M} in the free layer thus exerts a *spin torque*



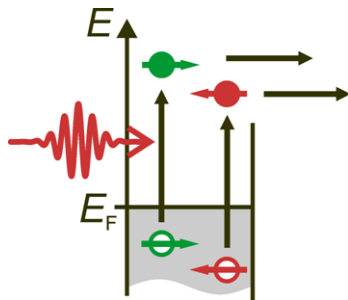
Spin-transfer by hot electrons



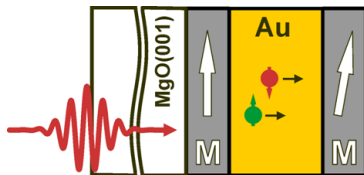
- a femtosecond laser can excite spin-polarized electrons as hot carriers



Spin-transfer by hot electrons



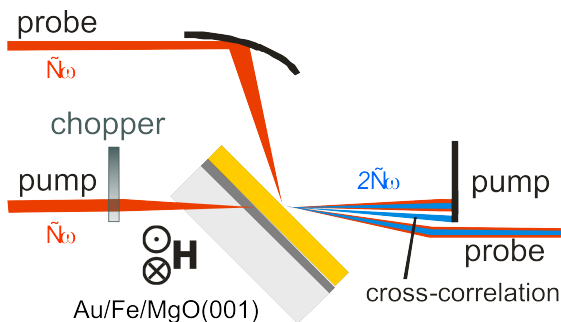
- a femtosecond laser can excite spin-polarized electrons as hot carriers



- so, can these spin-polarized hot electrons generate spin-transfer torque?



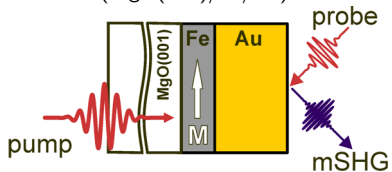
Optical pump-probe scheme



- cavity-dumped Ti:Sa oscillator, repetition rate: 76 MHz
- pulse-length: $\tau = 35$ fs
- energy: $E \approx 1.5$ eV \rightarrow Flux $\approx 1 \frac{\text{mJ}}{\text{cm}^2}$
- dumping rate: 1.5 MHz



Geometry of first samples
 (MgO(100)/Fe/Au):



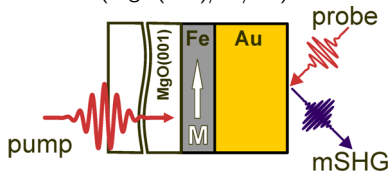
Time lapse of electron/spin-transport:
 (ballistic vs. diffusive):



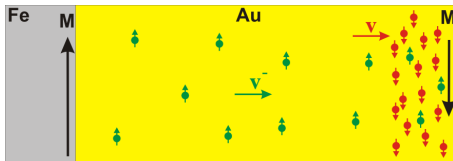
optical excitation of hot-carriers @ 0 fs



Geometry of first samples
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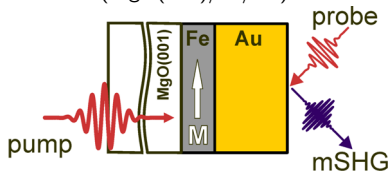
Time lapse of electron/spin-transport:
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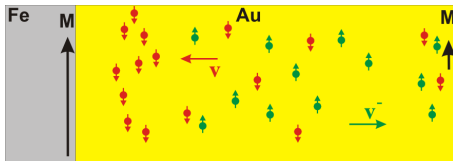
ballistic transport of minority carriers to surface @ 40 fs



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Time lapse of electron/spin-transport:
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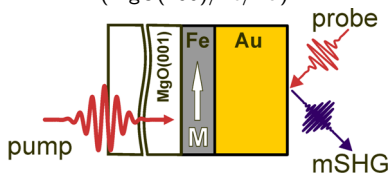


minority carriers are back-scattered at surface @ 150 fs

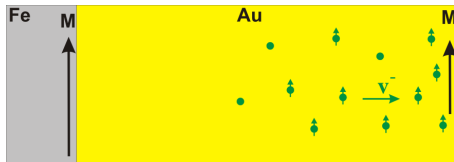
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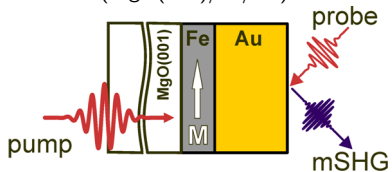


majority carriers reach surface diffusively @ 400 fs

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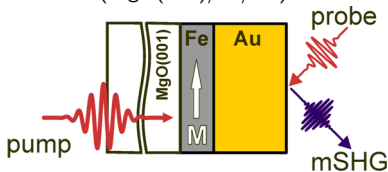


spin-disordering, thus magnetization relaxes @ 1000 fs

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Geometry of first samples
 (MgO(100)/Fe/Au):



Electric field:

$$\begin{aligned}\vec{E}_{2\omega} &= \vec{E}_{2\omega}^{\text{even}} + \vec{E}_{2\omega}^{\text{odd}} \\ &= \vec{\beta} + \vec{\alpha} M\end{aligned}$$

Intensity (measured signal):

$$I_{2\omega}^{\uparrow\downarrow} \propto |\vec{E}_{2\omega}^{\text{even}} + \vec{E}_{2\omega}^{\text{odd}}|^2 = |\vec{E}_{2\omega}^{\text{even}}|^2 \pm 2|\vec{E}_{2\omega}^{\text{even}} \vec{E}_{2\omega}^{\text{odd}}|$$

* the sign of the cross-term depends on direction of \vec{B}

Time lapse of electron/spin-transport:
 (ballistic vs. diffusive):

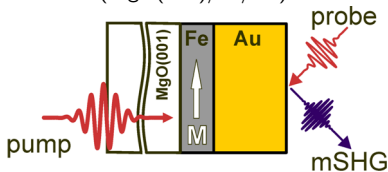


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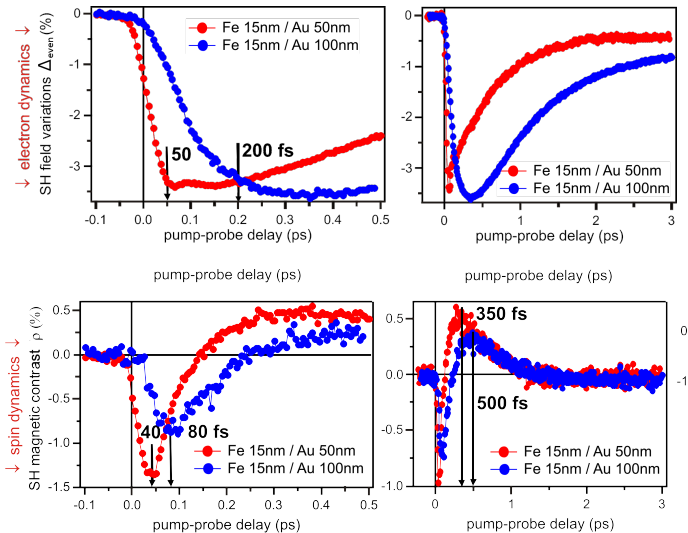
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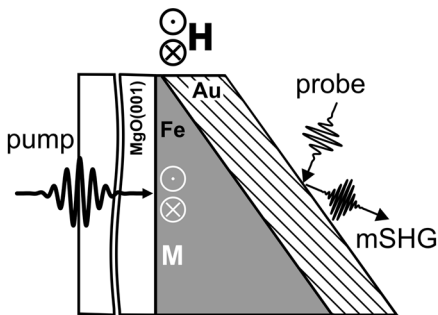
* the sign of the cross-term depends on direction of \vec{B}

We are interested in relative quantities:

- $\Delta_{\text{even}} \approx \frac{E_{2\omega}^{\text{even}}(t)}{E_0^{\text{even}}} - 1$
- $\Delta_{\text{odd}} \approx \frac{E_{2\omega}^{\text{odd}}(t)}{E_0^{\text{odd}}} - 1 \approx \frac{\Delta M(t)}{M_0}$
- $\rho(t) \approx 2 \frac{E_{2\omega}^{\text{odd}}(t)}{E_{2\omega}^{\text{even}}(t)} \propto M(t)$

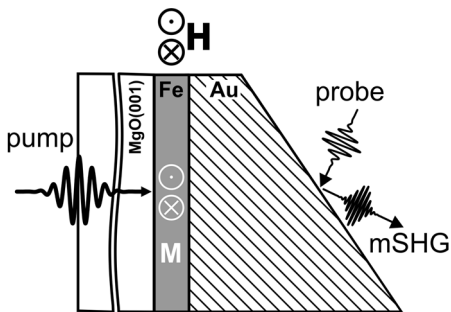






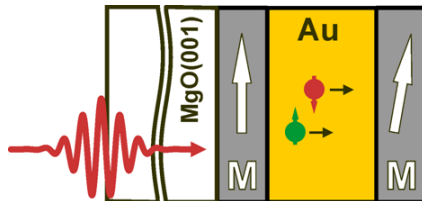
- to find the optimal thickness for the iron layer, we grow an iron wedge then scan along the wedge laterally
- optimal thickness lies between light penetration depth and ballistic length of the hot electrons





- grow a wedge analog to the iron wedge to find the proper thickness which suits both the ballistic mean free path of the electrons and the spin conservation length best





- Final Goal: After having determined optimal thicknesses for Fe and Au, grow a second Fe-layer onto Au with a free/different magnetization
- electrons from the first Fe-layer can now change magnetization in the second Fe-layer



Summary and Conclusion

- spin polarization occurs in ferromagnets due to band energy-splitting
- magnetization in thin films can be changed by direct spin-transfer with spin-polarized electrons
- spin-polarized carriers can be excited as hot carriers with a femtosecond laser
- minority carriers have higher velocities than majority carriers
- minority transport is mainly ballistically, majority mainly diffusively
- Final Goal: Spin transfer requires knowledge about optimal thicknesses for Fe/Au/Fe system

I would like to thank my working group and the DFG for funding our project.



