Solid State Theory Exercise 2

## Point groups and their representations

## Exercise 2.1 Energy bands of almost free electrons on the fcc lattice

Let us consider almost free electrons on a face-centered cubic (fcc) lattice. The goal of this exercise is to compute the lowest energy bands along the $\Delta$-line using degenerate perturbation theory and the machinery of the group theory. Remember that in reciprocal space, the fcc lattice transforms into a body-centered cubic (bcc) lattice. The point group of the cubic Bravais lattices (simple cubic, fcc, bcc) is denoted by $O_{h}$ (symmetry group of a cube). Its character table is given in Tab. 1.
a) We first study the $\Gamma$ point $(\vec{k}=0)$. For free electrons $(V=0)$ the lowest energy level is non-degenerate and the second one has a eight fold degeneracy. We focus on the second level and denote the eight-dimensional representation of $O_{h}$ defined on this subspace by $\Gamma$. Find the irreducible representations contained in $\Gamma$. Compute the group character $\chi_{\Gamma}$ and use the character table of $O_{h}$ to show that

$$
\begin{equation*}
\Gamma=\Gamma_{1}^{+} \oplus \Gamma_{2}^{-} \oplus \Gamma_{4}^{-} \oplus \Gamma_{5}^{+} \tag{1}
\end{equation*}
$$

| $O_{h}$ | $E$ <br> $[x y z]$ | $C_{3}(8)$ | $C_{4}^{2}(3)$ | $C_{2}(6)$ | $C_{4}(6)$ | $J$ | $J C_{3}(8)$ | $J C_{4}^{2}(3)$ | $J C_{2}(6)$ | $J C_{4}(6)$ |
| :---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: |
|  | $[z x y]$ | $[\bar{x} \bar{y} z]$ | $[y x \bar{z}]$ | $[\bar{y} x z]$ | $[\bar{x} \bar{y} \bar{z}]$ | $[\bar{z} \bar{x} \bar{y}]$ | $[x y \bar{z}]$ | $[\bar{y} \bar{x} z]$ | $[z \bar{x} \bar{z}]$ |  |
| $\chi_{\Gamma_{1}^{+}}$ | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 |
| $\chi_{\Gamma_{1}^{-}}$ | 1 | 1 | 1 | 1 | 1 | -1 | -1 | -1 | -1 | -1 |
| $\chi_{\Gamma_{2}^{+}}$ | 1 | 1 | 1 | -1 | -1 | 1 | 1 | 1 | -1 | -1 |
| $\chi_{\Gamma_{2}^{-}}$ | 1 | 1 | 1 | -1 | -1 | -1 | -1 | -1 | 1 | 1 |
| $\chi_{\Gamma_{3}^{+}}$ | 2 | -1 | 2 | 0 | 0 | 2 | -1 | 2 | 0 | 0 |
| $\chi_{\Gamma_{3}^{-}}$ | 2 | -1 | 2 | 0 | 0 | 2 | 1 | -2 | 0 | 0 |
| $\chi_{\Gamma_{4}^{+}}$ | 3 | 0 | -1 | -1 | 1 | 3 | 0 | -1 | -1 | 1 |
| $\chi_{\Gamma_{4}^{-}}$ | 3 | 0 | -1 | -1 | 1 | -3 | 0 | 1 | 1 | -1 |
| $\chi_{\Gamma_{5}^{+}}$ | 3 | 0 | -1 | 1 | -1 | 3 | 0 | -1 | 1 | -1 |
| $\chi_{\Gamma_{5}^{-}}$ | 3 | 0 | -1 | 1 | -1 | -3 | 0 | 1 | -1 | 1 |

Table 1: The character table of the cubic point group $O_{h}$.
b) A finite periodic potential will in general split the second energy level at the $\Gamma$ point. Applying degenerate perturbation theory to the Bloch equation [Eq. (1.20) in the lecture notes] leads to a $8 \times 8$ matrix with off-diagonal elements $u=V_{\frac{4 \pi}{a}(1,1,1)}$, $v=V_{\frac{4 \pi}{a}(1,0,0)}$ and $w=V_{\frac{4 \pi}{a}(1,1,0)}$ (we basically follow chapter 1.3 of the lecture notes). This matrix can be diagonalized by going into the symmetry subspaces.

Show that for the energies and the wave functions one finds

$$
\begin{array}{ccc}
\Gamma_{1}^{+}: & E_{0}+u+3 v+3 w & \cos \left(\frac{2 \pi}{a} x\right) \cos \left(\frac{2 \pi}{a} y\right) \cos \left(\frac{2 \pi}{a} z\right) ; \\
\Gamma_{2}^{-}: & E_{0}-u-3 v+3 w & \sin \left(\frac{2 \pi}{a} x\right) \sin \left(\frac{2 \pi}{a} y\right) \sin \left(\frac{2 \pi}{a} z\right) ; \\
\Gamma_{4}^{-}: & E_{0}-u+v-w & \left\{\sin \left(\frac{2 \pi}{a} x\right) \cos \left(\frac{2 \pi}{a} y\right) \cos \left(\frac{2 \pi}{a} z\right), \text { cyclic }\right\} ;  \tag{2}\\
\Gamma_{5}^{+}: & E_{0}+u-v-w & \left\{\cos \left(\frac{2 \pi}{a} x\right) \sin \left(\frac{2 \pi}{a} y\right) \sin \left(\frac{2 \pi}{a} z\right), \text { cyclic }\right\} ;
\end{array}
$$

where $E_{0}=\frac{\hbar^{2}}{2 m} 3\left(\frac{2 \pi}{a}\right)^{2}$.
c) How do the irreducible representations split on the $\Delta$-line? The $\Delta$-line is defined by the points $\vec{k}=\frac{\pi}{a}(0,0, \delta), 0 \leq \delta \leq 1$. Use the character table of $C_{4 v}$.

| $C_{4 v}$ | $E$ <br> $C_{2}(1)$ | $C_{4}(2)$ | $\sigma_{v}(2)$ | $\sigma_{d}(2)$ |  |
| :--- | ---: | ---: | ---: | ---: | ---: |
|  | $[x y z]$ | $[\bar{x} \bar{y} z]$ | $[y \bar{x} z]$ | $[\bar{x} y z]$ | $[y x z]$ |
| $\chi_{\Delta_{1}}$ | 1 | 1 | 1 | 1 | 1 |
| $\chi_{\Delta_{2}}$ | 1 | 1 | 1 | -1 | -1 |
| $\chi_{\Delta_{3}}$ | 1 | 1 | -1 | 1 | -1 |
| $\chi_{\Delta_{4}}$ | 1 | 1 | -1 | -1 | 1 |
| $\chi_{\Delta_{5}}$ | 2 | -2 | 0 | 0 | 0 |

Table 2: The character table of $C_{4 v}$.
d) Let us now consider the point $X=\frac{2 \pi}{a}(0,0,1)$. The lowest level is two fold and the second four fold degenerate for $V=0$. Compute the energies and the wave functions for these two levels.
e) Finally, sketch the energy bands between the $\Gamma$ and the $X$ point. For an actual numerical calculation use the values $u=-0.05, v=0.05$ and $w=0.1$ (in units of $\left.\frac{(2 \pi \hbar)^{2}}{2 m a^{2}}\right)$.

## Exercise 2.2 Lifting the degeneracy of the atomic states

Determine how the energy levels of the $p, d$ and $f$ orbitals of an atom lift due to a crystal field with cubic symmetry. Compute the corresponding eigenstates for the $d$ orbitals. For this, consider the homogeneous harmonic polynomials of order 2. Alternatively, have a look at the basis functions given on page 11 of the lecture notes.

