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## 2 Task Elektronenspektroskopie

## 2.1 (Film thickness)

This task is about a thin metal film on a substarte.

a)

We are meant to determine the film thickness, from the measured intensities of  $I_f$  and  $I_s$ . Therefore we first have to determine the intensities. We start with  $I_s$ , which is given with:

$$I_s = I_s^0 \int_0^\infty e^{-\frac{z}{\Lambda_{ss}}} dz \, e^{-\frac{d}{\Lambda_{ff}}}$$

while  $I_s^0 \propto n_s \sigma_s$ ,  $\Lambda_{ss} = \Lambda_s (E_s)$  and  $\Lambda_{ff} = \Lambda_f (E_f)$ . We use the substitution  $y = \frac{x}{\Lambda} \Leftrightarrow dx = \Lambda dy$  therefore we get:

$$I_s \propto n_s \sigma_s e^{-\frac{d}{\Lambda_{ff}}} \Lambda_{ss} \int_0^\infty dy \, e^{-y}$$

The integration is simple and only gives a 1:

$$I_s \propto n_s \sigma_s e^{-\frac{d}{\Lambda_{ff}}} \Lambda_{ss}$$

Now we can have a look at  $I_f$ , while this one is given with

$$I_f \propto n_f \sigma_f \int_0^d e^{-\frac{z}{\Lambda_{ff}}}$$

which leads to:

$$I_f \propto n_f \sigma_f \Lambda_{ff} \left( 1 - e^{-\frac{d}{\Lambda_{ff}}} \right)$$

Now we can get the quotient of the intensities to find d:

$$\frac{I_f}{I_s} = \frac{n_f \sigma_f \Lambda_{ff} \left(1 - e^{-\frac{a}{\Lambda_{ff}}}\right)}{n_s \sigma_s e^{-\frac{d}{\Lambda_{ff}}} \Lambda_{ss}}$$

which can be rearranged to:

$$d = \Lambda_{ff} \ln \left( 1 + \frac{n_s \sigma_s \Lambda_{ss} I_f}{n_f \sigma_f \Lambda_{ff} I_s} \right)$$

d therefore depends only on known quantities.

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To simplify this expression we assume, that the mean free path will be much bigger than the film thickness, meaning  $\Lambda \gg d$ . We can simplify the expression, starting with  $I_s$ 

$$I_s \propto n_s \sigma_s e^{-\frac{d}{\Lambda_{ff}}} \Lambda_{ss}$$

while we now have  $d \ll \Lambda_{ff}$  we can expand to first order:

$$I_s \propto n_s \sigma_s \Lambda_{ss}$$

and for  $I_f$ 

$$I_f \propto n_f \sigma_f \Lambda_{ff} \left( 1 - e^{-\frac{d}{\Lambda_{ff}}} \right)$$

while  $d \ll \Lambda_{ff}$ , we can expand the exponential function (to second order, while first order would cancel the whole term):

$$e^{-\frac{d}{\Lambda_{ff}}} \approx 1 - \frac{d}{\Lambda_{ff}} + \mathcal{O}\left(\left(\frac{d}{\Lambda_{ff}}\right)^2\right)$$

which leads to:

 $I_f \propto n_f \sigma_f d$ 

Now we can get the quotient of the intensities again in the simplified form to find d:

$$\frac{I_f}{I_s} = \frac{n_f \sigma_f d}{n_s \sigma_s \Lambda_{ss}}$$

which can again be rearranged to:

$$d = \frac{I_f n_s \sigma_s}{I_s n_f \sigma_f} \Lambda_{ss}$$

## 2.2 (Yield of photoelectrons in the high kinetic energy regime)

The law for the yield of photoelectrons in the high kinetic energy regime  $E_{kin} \approx \hbar \omega$  is given with  $Y \propto (\hbar \omega)^{\alpha}$ . We are meant to determine the exponent  $\alpha$ . We know that

$$Y \propto \int_0^z dz I$$
$$\propto \Lambda \propto E^{\frac{1}{2}}$$
$$\propto \sigma \propto E^{-3}$$

this means

 $Y \propto E^{-3 + \frac{1}{2}} \propto E^{\frac{5}{2}}$ 

or

 $Y\propto \sigma\cdot\Lambda$ 

We have a look at the figure called "Cross sections vs energy" in the lecture notes. We have three lines there, we will get an average value therefore:

line	$\sigma_e$
c1s	$-(2.9\pm 0.05)$
n1s	$-(3.0\pm0.05)$
o1s	$-(2.95\pm0.05)$
average	$-(2.95\pm0.03)$

Therefore we were able to determine the exponent with  $\sigma_e = -(2.95 \pm 0.03)$ . This first value seems to fit the expectation of -3 quite well, now the second graph leads to  $\Lambda_e = 0.5$ , which also fulfills our expectation. We therefore find:

$$Y \propto (\hbar\omega)^{\alpha} \propto (\hbar\omega)^{\sigma_e + \Lambda_e} \approx (\hbar\omega)^{-\frac{5}{2}}$$